The Effect of Pr co-doping on the Densification and Electrical Properties of Sr-LaAlO$_3$

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Perovskite-type materials have been investigated extensively in the attempt to discover new electrolyte materials for solid oxide fuel cells (SOFC) operating at low temperatures. These materials include LaAlO$_3$-based perovskite, which, when adequately doped, presents high ionic conductivity. However, despite the high conductivity of doped LaAlO$_3$, two characteristics limit its application as solid electrolyte: low sinterability and p-type electronic conductivity in oxidizing conditions. The purpose of this work was to investigate the effect of Pr co-doping on the densification and electrical conductivity of Sr-doped LaAlO$_3$. The addition of Pr eliminated the deficiencies mentioned above. Pr in solid solution acts as a perfect sintering aid because promoted densification without forming secondary phase which should be detrimental for electrical conduction. The addition of Pr also increased the bulk electrical conductivity measured in air atmosphere but had no effect at low oxygen partial pressure. However, the addition of Pr had a strong effect on the grain boundary, improving the ionic grain boundary electrical conductivity at air atmosphere which was attributed to the Pr influence on the space charge layer.

Keywords: SOFC, lanthanum aluminate, electrical characterization, solid electrolytes, space charge layer

1. Introduction

Solid oxide fuel cell (SOFC) technology is considered one of the most promising ways to generate electrical energy in the near future. This technology is based mainly on designs that employ yttrium-stabilized zirconia (YSZ) as electrolyte. The economic feasibility of fuel cells is directly related to the decrease in operating temperature, which is established mainly by the electrolyte’s properties$^{1,2}$. Thus, new electrolyte materials for SOFCs operating at low temperatures have been investigated extensively and perovskite-type materials are considered promising materials for electrolytes. These materials offer several advantages, including a stable crystal structure, a variety of elements that can be accommodated in the crystal lattice, and the relatively simple creation of oxygen vacancies by partial substitution of A- and/or B-site cations by lower valence cations$^{3-6}$.

In perovskite-based materials, anion migration occurs through the vacancy mechanism. To explain the ionic conductivity in perovskite-type oxides, the model of the ionic conduction mechanism involves structural parameters such as the tolerance factor, free volume and critical radius. The tolerance factor ($t$) described by Equation 1 establishes the relationship between the symmetry and ionic radius in perovskites. Most perovskites have a distorted cubic geometry. This distortion is associated with a high degree of anisotropy in oxygen sites. The free volume is defined as the “unoccupied volume” of the unit cell and is calculated by subtracting the volume occupied by all the ions in the

$$ t = \frac{r_A + r_O}{\sqrt{2}(r_B + r_O)} \quad (1) $$

where $r_A$ and $r_B$ represent the radiiuses of the cations and $r_O$ is the oxygen ion.

$$ r_{cv} = \frac{a_o((3/4)a_o - \sqrt{2}r_B) - (r_A - r_B)(r_A + r_B)}{2(r_A - r_B) + \sqrt{2}a_o} \quad (2) $$

where $a_o$ is the cubic lattice parameter.

Among the perovskite-structured oxides, LaGaO$_3$-based oxides have been investigated as electrolyte materials for intermediate-temperature SOFCs. Gallate-based materials show higher electrical conductivity than YSZ, but their chemical and mechanical stability are low at high temperatures. In fact, few materials can withstand both high temperature and low oxygen partial pressure, which is the condition to which the electrolyte is subjected on the anode region (e.g., T>800 °C and p(O$_2$)<10$^{-10}$ atm)$^{9-11}$. LaAlO$_3$-based materials are an alternative for LaGaO$_3$ compounds due to their stability in these conditions. LaAlO$_3$ presents high electrical conductivity when suitably doped. However, at high oxygen partial pressure p(O$_2$), the oxygen vacancy created by substituting La$^{3+}$ by cations of lower
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The electronic conductivity of doped LaAlO₃ was evaluated in air and at low oxygen partial pressure by the dc four-probe technique, using electron-blocking electrodes. Electrical measurements were taken with a HP 3468A multimeter in a temperature range of 250-800 °C. The electron-blocking electrodes were thin slices of YSZ with a thickness of <300 μm. One side of these YSZ slices was coated with platinum paste (Demetron 308A) and the other side was fixed onto the sample’s surface by mechanical pressure exerted by the sample holder inside the electrical measuring chamber. For measurements at low oxygen partial pressure, the samples were held for 1 hour at 800 °C and p(O₂)=10⁻² atm before taking the measurements.

3. Results and Discussion

3.1. Crystalline phases and densification

Figure 1 shows the X-ray diffraction patterns of sintered at 1600 °C/6h. The main diffraction peaks can be indexed based on the cubic perovskite structure of LaAlO₃ (JCPDS card 85-848), with La₆O₃ (JCPDS: 73-2141) detected as the minority phase.

Table 1 lists the lattice parameters, evaluated from XRD data, and the theoretical density which was calculated using the unit cell parameters determined by a Rietveld refinement of the X-ray diffraction pattern and considering 0.05 oxygen deficiency relative to oxygen sites. Refinements were carried out with the General Structure Analysis System (GSAS) software, using the graphical interface EXPGUI.

The crystal structures of the sintered samples were analyzed by X-ray diffraction (Siemens D5000 diffractometer) using Cu-Kα radiation (1.54 Å). The data were obtained in the 2θ range of 15-90° at a scan speed of 0.033°/s. The lattice parameters were evaluated from XRD data. Bulk densities of sintered specimens were measured by the Archimedes method. The relative density was determined from the theoretical value established from the experimental lattice parameters. The microstructures were analyzed by scanning electron microscopy (SEM) (Philips XL 30 FEG). The surfaces for SEM analysis were polished with 1-µm diamond paste and thermally etched at 100 °C lower than sintering temperature for 6 min.

The electrical conductivity of the specimens was measured in air and at low oxygen partial pressure by two-probe impedance spectroscopy technique (HP 4192 A) in a temperature range of 250 to 800 °C and a frequency range of 5 Hz to 13 MHz. Platinum paste (Demetron 308A) was applied on both sides of the pellets, which were then fired at 1100 °C for 30 min.

The sinterability was followed introducing different amount of Pr and sintering the samples at 1500 and 1600 °C. The electrical conductivity of sintered samples was evaluated by impedance spectroscopy in a broad range of temperatures and under high oxygen partial pressure p(O₂), to study the effect of Pr on the electrical conductivity of the LaAlO₃ based system. The electronic conduction in air was evaluated by the dc four-probe technique, using YSZ electron-blocking electrodes.

2. Experimental Procedure

2.1. Sample preparation

Sr-doped, nondoped and Pr co-doped LaAlO₃ powders were synthesized by the solid-state reaction method. A suitable amount of La₂O₃ (Aldrich, 99.9%), Al₂O₃ (Baikowski-CR30), Sr(NO₃)₂ (Riedel-de Haën, 99.99%) and Pr(NO₃)₃·4H₂O (Aldrich 99.9%) was milled with zirconia balls in isopropyl alcohol. The following nominal compositions were prepared: LaAlO₃, La₀₉₅Sr₀₅AlO₃, La₀₈₅Pr₀₁Sr₀₁AlO₃, and La₀₆₇Pr₀₂₁Sr₀₁AlO₃. The mixtures were calcined at 800 °C for 30 min and then milled into slurries in a vibratory mill with isopropyl alcohol and 1.0 wt% of polyvinyl butyl (B-98 Solutia). The slurries were dried in air and granulated through a nylon sieve. Pellets (10 mm diameter) were pressed isostatically (200 MPa) and sintered at 1500 and 1600 °C in air for a soak time of 6h.

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lattice parameter increased with the addition of Sr and decreased with the addition of Pr. Table 1 shows also the percentage of relative density of samples sintered at 1500 and 1600 °C where it can be seen the strong effect of Pr on the samples densification even at 1500 °C. Density as high as 97.6% of theoretical density was obtained for composition La$_{0.7}$(Pr$_{0.2}$Sr$_{0.1}$)AlO$_{3-δ}$ sintered at 1600 °C/6h.

Figure 2 shows the SEM micrographs obtained in back scattering electron (BSE) mode. The nondoped LaAlO$_3$ showed a homogeneous microstructure, while the Sr-doped samples presented dark gray grains dispersed in the continuous matrix of light gray grains. Since BSE imaging provides contrast due to differences in atomic numbers, dark gray grains are richer in an element with a low atomic number than light gray grains. SEM mapping of Pr undoped samples (Figure 3) showed that the dark gray grains were Al rich and La deficient. The addition of Pr improved the homogeneity of Al and La as show the SEM mapping of composition La$_{0.7}$(Pr$_{0.2}$Sr$_{0.1}$)AlO$_{3-δ}$, Figure 4. The addition of Pr, besides improved the densification, inhibited the grain growth and the La deficient grains, dark gray grains, are smaller than matrix grains (Figure 2). The average grain size of La$_{0.8}$Sr$_{0.1}$AlO$_{3-δ}$, La$_{0.9}$(Pr$_{0.1}$Sr$_{0.1}$)AlO$_{3-δ}$ and La$_{0.7}$(Pr$_{0.2}$Sr$_{0.1}$)AlO$_{3-δ}$ were 1.83 ± 0.08 µm, 1.44 ± 0.08 µm and 1.1 ± 0.2 µm, respectively.

3.2. Electrical conductivity

The impedance spectroscopy spectra in Figure 5, which were measured in air at 375 °C, show two distinct semicircles. The high frequency semicircle is attributed to the bulk response and the low frequency to the grain

Table 1. Lattice parameter and theoretical and relative densities of sintered samples.

<table>
<thead>
<tr>
<th>Composition</th>
<th>Lattice parameter (Å)</th>
<th>Theoretical density (g/cm$^3$)</th>
<th>%Relative Density 1500 °C</th>
<th>%Relative Density 1600 °C</th>
</tr>
</thead>
<tbody>
<tr>
<td>LaAlO$_3$</td>
<td>3.790</td>
<td>6.52</td>
<td>99.2</td>
<td>99.8</td>
</tr>
<tr>
<td>La$<em>{0.9}$Sr$</em>{0.1}$AlO$_{3-δ}$</td>
<td>3.793</td>
<td>6.33</td>
<td>77.0</td>
<td>90.1</td>
</tr>
<tr>
<td>La$<em>{0.8}$(Pr$</em>{0.1}$Sr$<em>{0.1}$)AlO$</em>{3-δ}$</td>
<td>3.786</td>
<td>6.37</td>
<td>93.6</td>
<td>94.7</td>
</tr>
<tr>
<td>La$<em>{0.7}$(Pr$</em>{0.2}$Sr$<em>{0.1}$)AlO$</em>{3-δ}$</td>
<td>3.789</td>
<td>6.46</td>
<td>93.0</td>
<td>97.6</td>
</tr>
</tbody>
</table>

Figure 2. SEM micrograph of polished and thermally etched surfaces of samples sintered at 1600 °C/6h.
boundary. However, the higher frequency semicircle disappeared gradually with increasing temperature.

3.2.1. Influence of Pr-doping on the bulk and grain boundary electrical conductivity

Figure 6 depicts Arrhenius plots of the bulk electrical conductivity measured in air which increases systematically with increasing concentration of Pr.

As mentioned earlier, ionic conductivity depends on structural parameters. Oxide ions can move easily to neighboring oxygen vacancy sites when the tolerance factor is near unity, i.e., the deformation of the cubic symmetry is small. Thus, the higher the symmetry of the structure the higher is the conductivity of oxide ions. A large free volume favors high oxygen ion mobility, and high values of critical radius result in low disturbances of the surrounding cations during the migration of the anion through the lattice. Considering that the valence of Pr is equal to or higher than that of La, its presence in crystal lattice does not introduce additional oxygen vacancies. Therefore, the high conductivity of the Pr co-doped samples can be attributed to enhancement of oxygen ion mobility due to a more favorable crystalline structure. In other words, the addition of Pr increased systematically the critical radius and the free volume and reduced the distortion of the crystal lattice (Table 2). The decrease in the bulk activation energy in response to the addition of Pr (see Table 3) reinforces the notion of a correlation between the crystalline structure and the high oxygen ion mobility of the Pr co-doped samples.

Figure 7 shows Arrhenius plots for microscopic grain boundary conductivity measured in air. The microscopic grain boundary conductivity ($\sigma_{gb}^{\text{micro}}$) can be calculated by Equation 4:

$$\sigma_{gb}^{\text{micro}} = \frac{\varepsilon_o \varepsilon_{gb}}{R_{gb} C_{gb}}$$

where $\varepsilon_o$ is the vacuum permittivity, and $\varepsilon_{gb}$, $R_{gb}$ and $C_{gb}$ are the grain boundary permittivity, resistance and capacitance, respectively. Assuming that the grain boundary permittivity $\varepsilon_{gb}$ is equal to the bulk permittivity, $\varepsilon_b$, can be calculated by Equation 5:
$\varepsilon_b = \frac{C_b L}{\varepsilon_0 A}$

$\sigma_{gb}^{micro} = \frac{C_b L}{R_{gb} C_{gb} A}$

where $L$ and $A$ are, respectively, the sample’s thickness and electrode area.

The $\sigma_{gb}^{micro}$ shown in Figure 7 was calculated by Equation 6:

The addition of Pr improved systematically the microscopic grain boundary conductivity (see Figure 7) and decreased the grain boundary activation energy (Table 3). Since additional vacancies were not introduced with the
addition of Pr, the behavior observed here can be attributed to the influence of Pr on the grain boundary thickness ($\delta_{gb}$), as has been exhaustively discussed for doped-ceria and yttria-stabilized zirconia\textsuperscript{[4]}. The $\delta_{gb}$ values for Pr-doped and nondoped samples analyzed at 425 °C in air were calculated by Equation 7 and are shown in Table 3.

![Figure 5. Impedance spectroscopy spectra, measured in air at 375 °C, of the samples sintered at 1600 °C/6h.](image)

\[ \delta_{gb} = \frac{\varepsilon_{gb}}{\varepsilon_g C_{gb} d_g} \]  
(7)

where $\varepsilon_g$ and $\varepsilon_{gb}$ are the grain boundary and grain permittivity that were considered equal, $C_g$ and $C_{gb}$ are, respectively, the capacitance for grain and grain boundary, and $d_g$ is the average grain size. The systematically decrease of $\delta_{gb}$ with increasing of Pr amount, as shown Table 3, confirms the positive effect of Pr doping on the decrease of the grain boundary thickness.

Although measuring the electrical conductivity by impedance spectroscopy enables the separation of bulk and grain boundary conductivity, it did not provide information about the ionic and electronic contributions to the total electrical conductivity. Since the composition $\text{La}_{0.87}\text{(Pr}_{0.1}\text{Sr}_{0.1})\text{AlO}_3$ showed the highest electrical conductivity in air it was analyzed by the dc four-probe technique using YSZ electron-blocking electrodes\textsuperscript{34}. For comparison, the Pr nondoped composition, $\text{La}_{0.87}\text{Sr}_{0.1}\text{AlO}_3$, was also analyzed by this technique. Figure 8 compares the Arrhenius plots for the dc electrical conductivity of the $\text{La}_{0.87}\text{Sr}_{0.1}\text{AlO}_3$ and $\text{La}_{0.87}\text{(Pr}_{0.1}\text{Sr}_{0.1})\text{AlO}_3$ samples obtained in air with Pt and YSZ electron-blocking electrodes. As can be seen, the electrical conductivity measured with electron-blocking electrodes (ionic conduction) was lower than that measured with Pt electrodes (total conduction) because the oxygen vacancies are filled in an air atmosphere, generating electron holes, according to equation 3. The following was observed: $i)$ the ionic conductivity of $\text{La}_{0.87}\text{(Pr}_{0.1}\text{Sr}_{0.1})\text{AlO}_3$ is higher than that of $\text{La}_{0.87}\text{Sr}_{0.1}\text{AlO}_3$, and $ii)$ the decrease in electrical conductivity measured with the electron-blocking electrode of the Pr-doped sample was lower (6 times) than the Pr-nondoped sample (8 times). This indicates that no additional oxygen vacancies were introduced by the addition of Pr, reinforcing the hypothesis that the addition of Pr enhances oxygen ion mobility due to a more favorable crystalline structure, resulting in high electrical conductivity. The total electrical conductivity of the Pr-doped sample at 800 °C was higher than Pr-nondoped sample, 23.0 x 10\textsuperscript{-3} S/cm and 9.0 x 10\textsuperscript{-3} S/cm, respectively. The value obtained for Pr-nondoped in the literature is 6.3 x 10\textsuperscript{-3} S/cm\textsuperscript{34}.

3.2.2. Influence of oxygen partial pressure on the bulk and grain boundary electrical conductivity

Figure 9 compares the Arrhenius plots of the bulk and microscopic grain boundary electrical conductivity of the $\text{La}_{0.87}\text{Sr}_{0.1}\text{AlO}_3$, $\text{La}_{0.87}\text{(Pr}_{0.1}\text{Sr}_{0.1})\text{AlO}_3$, and $\text{La}_{0.87}\text{(Pr}_{0.1}\text{Sr}_{0.1})\text{AlO}_3$.

<table>
<thead>
<tr>
<th>Composition</th>
<th>Critical Radius (Å)</th>
<th>Free Volume (Å\textsuperscript{3})</th>
<th>Tolerance Factor</th>
</tr>
</thead>
<tbody>
<tr>
<td>$\text{La}<em>{0.87}\text{Sr}</em>{0.1}\text{AlO}_3$</td>
<td>0.9016</td>
<td>14.06</td>
<td>1.011</td>
</tr>
<tr>
<td>$\text{La}<em>{0.87}\text{(Pr}</em>{0.1}\text{Sr}_{0.1})\text{AlO}_3$</td>
<td>0.9128</td>
<td>14.16</td>
<td>1.003</td>
</tr>
<tr>
<td>$\text{La}<em>{0.87}\text{(Pr}</em>{0.1}\text{Sr}_{0.1})\text{AlO}_3$</td>
<td>0.9296</td>
<td>14.44</td>
<td>0.994</td>
</tr>
</tbody>
</table>

3.2.2. Influence of oxygen partial pressure on the bulk and grain boundary electrical conductivity

Table 3. Activation energies for the bulk ($E_{ab}$) and grain boundary ($E_{ab}$) and grain boundary thickness ($\delta_{gb}$) at 425 °C (nm) of samples measured in air and at low oxygen partial pressure.

<table>
<thead>
<tr>
<th>Composition</th>
<th>$E_{ab}$ (eV)</th>
<th>$E_{ab}$ (eV)</th>
<th>$\delta_{gb}$ at 425 °C (nm)</th>
</tr>
</thead>
<tbody>
<tr>
<td>$\text{La}<em>{0.87}\text{Sr}</em>{0.1}\text{AlO}_3$</td>
<td>In air, $p(O_2) = 10^{-22}$</td>
<td>0.96</td>
<td>1.81</td>
</tr>
<tr>
<td>$\text{La}<em>{0.87}\text{(Pr}</em>{0.1}\text{Sr}_{0.1})\text{AlO}_3$</td>
<td>In air, $p(O_2) = 10^{-22}$</td>
<td>0.85</td>
<td>1.24</td>
</tr>
<tr>
<td>$\text{La}<em>{0.87}\text{(Pr}</em>{0.1}\text{Sr}_{0.1})\text{AlO}_3$</td>
<td>In air, $p(O_2) = 10^{-22}$</td>
<td>0.88</td>
<td>1.64</td>
</tr>
</tbody>
</table>

Figure 6. Arrhenius plot of bulk electrical conductivity measured in air for Pr doped and nondoped samples sintered at 1600 °C/6h.

Figure 7. Arrhenius plot of grain boundary (gb) electrical conductivity measured in air for Pr doped and nondoped samples sintered at 1600 °C/6h.

Figure 8. Arrhenius plots of the bulk and grain boundary electrical conductivity measured in air for Pr doped and nondoped samples sintered at 1600 °C/6h.
AlO$_{3.8}$ compositions measured at a low oxygen partial pressure, P$_{O2}$=10$^{-12}$ atm and in air atmosphere.

All compositions showed lower bulk conductivity at low oxygen partial pressure than in air atmosphere, but the Pr-nondoped sample presented a lower decrease than the La$_{0.7}$(Pr$_{0.2}$Sr$_{0.1}$)AlO$_{3.8}$, which was the most conductive sample. In addition, the bulk activation energies at low oxygen partial pressure were only ~3% higher than in air atmosphere (Table 3). This small difference means that the bulk characteristic did not change significantly at low P$_{O2}$.

However, the microscopic grain boundary conductivities ($\sigma_{gb}^{micro}$) of all compositions were highly sensitive to the oxygen partial pressure. The $\sigma_{gb}^{micro}$ at low P$_{O2}$ was about two orders of magnitude lower than in air (Figure 9), while the grain boundary activation energy increased (Table 3). This behavior can be understood by analyzing the effect of P$_{O2}$ on the grain boundary thickness, $\delta_{gb}$, whose values at 425 °C are listed in Table 3. Note that $\delta_{gb}$ showed a significant increase with decreasing P$_{O2}$, which is consistent with the decrease in $\sigma_{gb}^{micro}$ and the increase in activation energy (Table 3).

Although Pr-doped and nondoped samples displayed the same behavior in response to variations in oxygen partial pressure, the lower increase in grain boundary activation energy and $\delta_{gb}$ with the decrease in P$_{O2}$ presented by the Pr-doped samples provided evidence that Pr acted positively on the grain boundary.

The lower decrease in bulk electrical conductivity of the Pr-nondoped sample under low oxygen partial pressure (Figure 9) can be attributed to the high grain boundary thickness formed in this sample under low oxygen partial pressure (Table 3), which inhibited oxygen diffusion from the bulk, thus preventing the elimination of electronic conduction.

4. Conclusions

Praseodymium as a co-dopant in Sr-LaAlO$_3$ behaves as an excellent sintering aid besides increasing the bulk electrical conductivity measured in air atmosphere. Oxygen ion mobility is enhanced by the favorable crystalline structure established by the Pr solid solution. At low P$_{O2}$, the Pr co-doping has no effect on bulk conductivity.

The Pr co-doping has a strong effect on the grain boundary, improving the ionic conductivity due to change in the space charge layer structure. At high oxygen partial pressure, the grain boundary thickness ($\delta_{gb}$) of Pr doped samples (3.8 and 4.4 nm) was lower than that of samples...
without Pr (5.7 nm). The same behavior was observed at a low oxygen partial pressure, in addition to the significant increase in space charge thickness. These results indicate that Pr improved the total electrical conductivity through its positive effect on the grain boundary.

For Pr-nondoped samples, the small difference between bulk conductivity measured at low and high $P_{O_2}$ can be attributed to the high grain boundary thickness, which inhibits oxygen diffusion from the bulk, thereby preventing the total elimination of the electronic conduction.

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