

Vibrational Spectra of Bis(dmit) Complexes of Main Group Metals: IR, Raman and *Ab Initio* Calculations

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Este trabalho apresenta uma investigação teórico-experimental dos espectros infravermelho/Raman de diversos metais representativos coordenados com 1,3-ditiole-2-tiona-4,5-ditioato (dmit), a saber [NEt₄]₂[Zn(dmit)₂], [NEt₄][Sb(dmit)₂] e [NEt₄][Bi(dmit)₂]. Os espectros IV/Raman de todos os sólidos e também o espectro IV em solução do composto [NEt₄]₂[Zn(dmit)₂] foram registrados de 4000 a 100 cm⁻¹. Duas regiões foram claramente identificadas: abaixo de 380 cm⁻¹ os modos apresentaram forte contribuição metal-ligante, e acima, os modos indicaram majoritária contribuição do ligante. Os espectros IV/Raman do NEt₄Br também foram registrados. As bandas vibracionais foram atribuídas a partir da análise de compostos relacionados, de publicações prévias e comparadas com amplo conjunto de cálculos *ab initio* realizados com diversos ECPs (SBK, Stuttgart), conjunto de bases (com e sem funções d) e metodologias (RHF, DFT e MP2). As geometrias, frequências e intensidades foram calculadas e comparadas com as frequências experimentais visando atribuições seguras. Assim, uma excelente concordância foi obtida para diversas bandas experimentais.

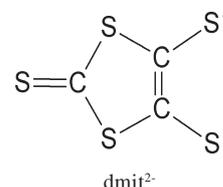
This work reports a theoretical-experimental investigation of the infrared/Raman vibrational spectra of several metal 1,3-dithiole-2-thione-4,5-dithiolate (dmit) complexes; [NEt₄]₂[Zn(dmit)₂], [NEt₄][Sb(dmit)₂] and [NEt₄][Bi(dmit)₂]. IR/Raman spectra of all the solids and the solution IR spectrum of [NEt₄]₂[Zn(dmit)₂] were recorded from 4000 to 100 cm⁻¹. Two regions were clearly identified: below 380 cm⁻¹, the modes presented significant metal-ligand contributions, and above, the modes indicated major ligand contributions. IR/Raman spectra of NEt₄Br have also been recorded. The vibrational bands have been assigned starting from the analysis of related compounds, of previous published spectra and compared to several ECPs (SBK, Stuttgart), basis sets (with and without functions d) and methodologies (RHF, MP2 and DFT). Geometries, frequencies and intensities were calculated and compared to the experimental frequencies seeking secure assignment. Thus, an excellent agreement was obtained for several experimental bands.

Keywords: IR, Raman, *ab initio*, RHF, DFT, MP2, dmit complexes, Zn, Sb, Bi

Introduction

The syntheses and properties of metal complexes with 1,2-dithiolato-polysulfido ligands have been intensely studied¹⁻³ particularly since the discovery, by Cassoux and co-workers⁴ in 1986, of the first metallic molecular superconductor, [TTF][Ni(dmit)₂], where dmit is the 1,3-dithiole-2-thione-4,5-dithiolate ligand. Presently dmit ligand complexes are the best studied.

The publication of Cassoux and co-workers⁴ seminal paper led a new impetus for the study of the applications



of metal dmit compounds as super-conducting complexes, materials with non-linear optical properties and materials for electronic molecular devices. Such properties arise in part, from the rich redox chemistry of these compounds and the large polarizability of sulfur atoms.

An additional important feature for the super-conducting property on metal bis(dmit) complexes, such as [TTF][Ni(dmit)₂], is the planar geometry of the anions,³⁻⁵

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which allows the formation of one and three-dimensional packing stacks in the solid state. However, a wide range of symmetries, coordination shells and solid state arrangements have been reported for $[M(\text{dmit})_2]^{y/x}$ in ionic $[Q]_x[M(\text{dmit})_2]^{y/x}$ species [e.g., $Q = \text{onium cation}$].¹⁻¹⁵ For instance, D_{2d} symmetry is present in Cu(II) and Zn(II) complexes, D_{2h} symmetry in Ni(II) complexes and C_2 symmetry for complexes of the Group 15 elements, Sb and Bi, in their trivalent oxidation state.

Despite the large volume of work on metal-dmit complexes, very few efforts have been devoted to the interpretation of their spectroscopic features. Actually most of the reported IR and Raman spectra have been used essentially for identification and characterization purposes.¹⁸⁻²⁴ A full vibrational assignment of the IR and Raman spectra for most of the metal-dmit complexes is now required to assess features of the charge-transfer process and the superconductivity properties. Therefore we are presently starting a careful analysis of IR and Raman spectra concerning fundamental vibrational assignment, overtones, combination bands, and extensive *ab initio* calculations of all such properties. Such information is considered to be essential in providing a clear understanding of properties such as local symmetry and the crystal environment, inter and intra-molecular charge-transfer processes and its correlation to the molecular and crystalline vibration spectra.

The main goal of the work now reported was to study the vibrational spectra of the representative metal bis(dmit) complexes, $[\text{NEt}_4]_2[\text{Zn}(\text{dmit})_2]$, $[\text{NEt}_4][\text{Sb}(\text{dmit})_2]$ and $[\text{NEt}_4][\text{Bi}(\text{dmit})_2]$, and to provide a full assignment of their IR and Raman bands, assisted by *ab initio* calculations carried out with different effective core potentials, basis sets and calculation levels. Up to now, the only species previously studied with *ab initio* calculations and Raman spectroscopy were $\text{TTF}[\text{Ni}(\text{dmit})_2]$, $\text{TTF}[\text{Pd}(\text{dmit})_2]$ and the neutral $[\text{M}(\text{dmit})_2]$ ($M = \text{Ni, Pd, Pt}$).²⁰⁻²² A low resolution Raman spectrum of $[\text{NEt}_4]_2[\text{Zn}(\text{dmit})_2]$ has been previously reported;²⁴ however it suffered from sample decomposition in the laser beam.

The choice of bis(1,3-dithiole-2-thione-4,5-dithiolato)-zincate, -antimonate and -bismuthate complexes presents a very convenient probe for family behaviour study. All the crystal structures of these complexes are known,⁶⁻¹⁵ including those of the tetraethylammonium salts, a family showing an increasing degree of inter-anion interactions from the Zn to Bi complex. Therefore we could assess M-S inter-anion interactions in the solid state, by carefully comparing the gas-phase calculation to the solid-state bulk spectra of these systems. The $[\text{Zn}(\text{dmit})_2]^{-2}$ anions in complex salts $[Q]_2[\text{Zn}(\text{dmit})_2]$ ($Q = \text{NEt}_4, \text{NBU}_4, \text{ferrocenyl-CH}_2\text{NMe}_3, 1,4\text{-Me}_2\text{-pyridinium, PPh}_4 \text{ and AsPh}_4$)⁶⁻⁹ in the solid state, invariably have distorted tetrahedral

geometries, and four coordinate zinc centres, with no anion-anion interactions apart for very weak S---S contacts close to its van der Waals radii sum. The distortion from an ideal zinc tetrahedral coordination is a consequence of the near 95° bite angles of the chelating dmit ligands. Each dmit ligand is symmetrically coordinated, or nearly so, to zinc, with the Zn-S bond lengths for all compounds falling in the narrow range of 2.3285(8) to 2.3522(10) Å. While the cation, Q, has little or no effect on the zinc coordination geometry, the overall shape of the anion, and also the extent of the S---S contacts, do vary with Q.

The situation is different in the case of the anions in $[Q][M(\text{dmit})_2]$ ($M = \text{Sb}^{10-12}$ or Bi^{13-15}): here there are metal-sulfur inter-anion interactions, that raise the coordination number of the metal to 6, and that including the lone pair on the metal, provide highly distorted pseudo pentagonal bipyramidal geometries for Sb and Bi. While the cation, Q^+ , in each of the series, influences the anion arrangements and also the very weak S---S inter-anion interactions, the basic geometries about each of the two metals remain unaffected. The inter-anion M-S interactions, either with thiolato or thione S atoms, are considerably weaker than the intra-anion M-S bonds. As each of the dmit ligands is asymmetrically bound to the metal, three distinct regions of M-S bond lengths are found. For $[Q][\text{Sb}(\text{dmit})_2]$ ($Q = \text{NEt}_4, \text{PPh}_4 \text{ and AsPh}_4$), the Sb-S bond lengths fall into the ranges: 2.4679(18) to 2.527(3), 2.6047(15) to 2.7311(6) [both intra anion bonds] and 3.271(4) to 3.5495(15) Å [inter-anion bond]. For $[Q][\text{Bi}(\text{dmit})_2]$ ($Q = \text{NEt}_4, \text{NBU}_4 \text{ and Ph}_4\text{As}$), the Bi-S bond lengths fall into three groups: 2.5872(10) to 2.527(3), 2.7299(11) to 2.87(1) [both intra bonds] and 3.0391(14) to 3.4628(13) Å [inter bonds]. Comparisons of the bond lengths suggest that the inter-anion interactions are relatively stronger for Bi than for Sb.

Experimental

Compounds

The complexes $[\text{NEt}_4]_2[\text{Zn}(\text{dmit})_2]$,⁷ $[\text{NEt}_4][\text{Sb}(\text{dmit})_2]$ ¹² and $[\text{NEt}_4][\text{Bi}(\text{dmit})_2]$ ¹⁵ were obtained through well known published procedures. All complexes had physical properties, melting point and elemental analysis in agreement with the published values. A commercial sample of tetraethylammonium bromide was used as received.

Ab initio calculations

All calculations here presented have been designed to explore vibrational assignment of synthesized inorganic materials. Therefore the qualities sought in the calculations

are robustness, reliability, facility to use, and economical, *i.e.*, it must be executed on personal computers presently available. Therefore, it should employ ECPs with double-zeta basis function set easily available with RHF, DFT or, at most, MP2 methodologies.

Previous experiences showed that minimal basis provide unreliable frequencies and assignments while triple-zeta plus polarization led to intense geometry corrections but minor frequency switches when compared to our results. Considering the excessive costs and minor modifications we shall explore a compact easy-to-use ECP + double-zeta + polarization scheme along this work.

The calculations employed the programs GAMESS R5²⁵ on a Red Hat 7.3 Linux cluster and the Gaussian 98²⁶ on a Window platform. The 6-31G basis set was employed for all elements except the central metallic atom. This basis set was chosen due to our previous experience with this basis set in several vibrational studies carried out in this group.^{33,34} This experience shows most frequencies regularly scaled by a factor of 0.8953.³³ For the central metal atom, different basis set were tested. In the SBK calculation, the compact effective pseudopotential of Stevens *et al.*²⁷ with double-zeta basis set was employed for the metal atom. We also applied this basis set extended with metal d functions (SBKd). In the Stuttgart²⁸ pseudopotential/basis set, the electron core potential (ECP) from Stuttgart University was used.

Three alternative methods were studied: restricted Hartree-Fock calculations (RHF), Moller-Plessett perturbation theory of second order (MP2), and the functional density theory (DFT-B3LYP) by employing the B3 potential for the coulomb and LYP for the exchange correlation term. All calculations started with SCF convergence, geometry optimisation and after obtaining a stable minimum, the frequencies were calculated. This minimum was accepted as long as no imaginary molecular frequencies were detected. Once the normal modes were obtained, the IR and Raman intensities were evaluated. All modes were visually characterized through the graphical MOLDEN 3.7 and Molekel programs.^{29,30}

These calculations provide qualitative and quantitative consistent results for each of the three complexes. For this reason, only tables and figures related to the SBKd ECP/basis set will be presented. However, whenever a significant difference is observed between the different ECPs, this will be indicated in the text.

Vibrational spectra

The infrared spectra, from 4000 to 150 cm⁻¹, were recorded in CsI pellets using a Nicolet Magna-IR 760

Fourier transform instrument at room temperature, with a 4 cm⁻¹ resolution. The FT-IR spectrum of [NEt₄]₂[Zn(dmit)₂], from 4000 to 600 cm⁻¹, was also recorded in dichloromethane solutions [concentrations: 5, 2 and 0.5 m/v].

Raman spectra were obtained as powder in glass capillaries on a Nicolet FT Raman 950 spectrophotometer. The spectra were recorded at room temperature with a germanium detector, maintained at liquid nitrogen temperature and using 1064.0 nm radiation, generated by a Nd-YAG laser with a resolution of 2 cm⁻¹. The samples were exposed to an average of 500 to 800 mW of laser intensity during the spectral recording.

Results

[Zn(dmit)₂]⁻² geometry optimization

The [Zn(dmit)₂]⁻² anion, of D_{2d} symmetry, has the highest symmetry found among the studied anions. The absence of significant interactions with neighbour dmit units should result in reduced environment effects in the IR and Raman spectra.

Figure 1 presents the [Zn(dmit)₂]⁻² optimized structure and its atomic numbering scheme. Table 1 lists the main geometrical parameters obtained from the geometry optimization with the basis SBKd at RHF, DFT and MP2 level: no significant differences were found between the gas phase calculated values and those obtained in the crystal structure determination. Calculated parameters using the Stuttgart effective core potential were similar and hence are not displayed here. As the ligands share the same basis sets, any difference between these calculation procedures will result from the metal ECP and its basis set. Zn(II) has a filled d shell and provides its unoccupied 4s and 4p orbitals for binding to the ligands. Therefore the most important electrons in the Zn-S bonding must arise from the ligands.

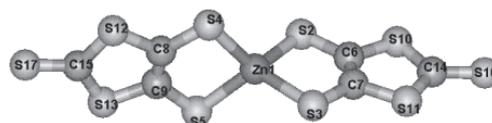


Figure 1. Three-dimensional representation of complex anion [Zn(dmit)₂]⁻² calculated with ab initio method.

Calculations using the MP2 methodology provide Zn-S bond lengths closest to those determined experimentally. Furthermore, intra-ligand bond lengths increase and the Zn-S bond length decrease in going from the RHF to the MP2 and DFT calculations. This might be an indication that the main correlation energy process is related to the electronic delocalization from the dmit ligands toward the central metallic atom.

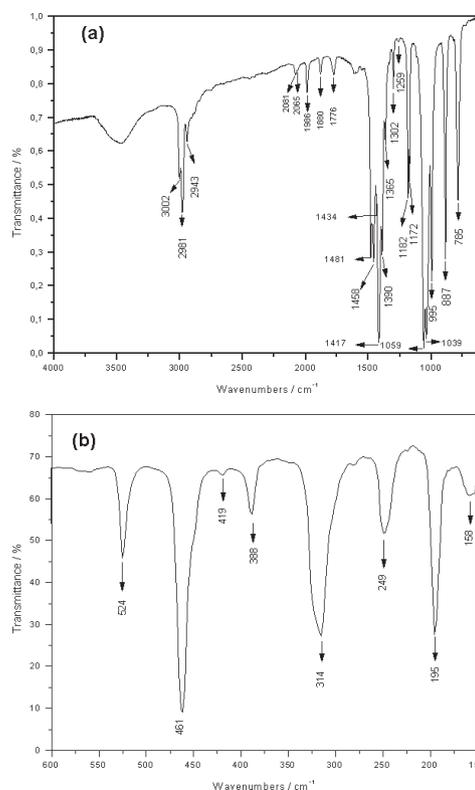
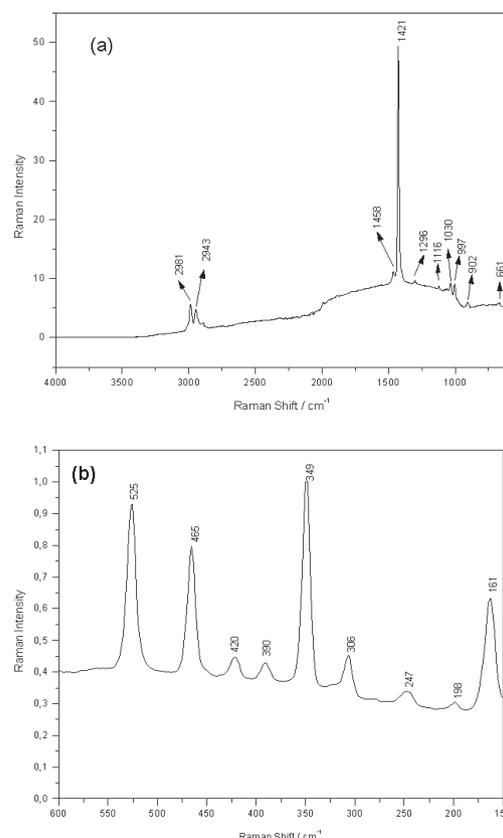
Table 1. Main geometrical parameters for the complex anion $[\text{Zn}(\text{dmit})_2]^{2-}$ with ab initio calculations and the SBKd basis set

	RHF	DFT	MP2	Exp. ⁷
<i>Bond (Å)</i>				
S ¹⁶ =C ¹⁴	1.7055	1.7154	1.7274	1.6590
S ¹¹ -C ¹⁴	1.7706	1.7995	1.8057	1.7210
S ¹¹ -C ⁷	1.8338	1.8499	1.8475	1.7520
S ³ -C ⁷	1.8064	1.8035	1.8200	1.7415
C ⁶ =C ⁷	1.3285	1.3607	1.3730	1.3540
Zn-S ³	2.4473	2.4251	2.3903	2.3503
<i>Bending (°)</i>				
S ¹¹ C ¹⁴ S ¹⁶	123.8	124.2	124.0	123.7
C ⁷ S ¹¹ C ¹⁴	97.7	98.3	98.1	98.4
S ¹¹ C ⁷ C ⁶	116.1	115.8	115.9	115.3
ZnS ³ C ⁷	96.1	95.6	95.2	95.9
S ² ZnS ³	118.6	117.7	116.6	115.9
S ² ZnS ⁵	118.6	117.7	116.6	116.7
S ² ZnS ⁴	118.6	117.7	116.6	119.4
S ³ ZnS ⁴	118.6	117.7	116.6	113.1
<i>Dihedral (°)</i>				
ZnS ³ C ⁷ C ⁶	0.0	0.0	0.0	3.3
S ³ C ⁷ S ¹¹ C ¹⁴	180.0	180.0	180.0	177.6
C ⁷ S ¹¹ C ¹⁴ S ¹⁶	180.0	180.0	180.0	179.4

[Zn(dmit)₂]²⁻ experimental frequencies

All the IR/Raman peaks associated with the $[\text{NEt}_4]^+$ cation were found in the spectra of the complex, $[\text{NEt}_4]_2[\text{Zn}(\text{dmit})_2]$, see electronic supplementary information. The anion, $[\text{Zn}(\text{dmit})_2]^{2-}$, belongs to the D_{2d} point group. By analyzing the molecular cartesian displacement representation and its irreducible composition, all genuine vibrations are given by: $\Gamma_{\text{vibration}} = 8A_1 + 2A_2 + 3B_1 + 8B_2 + 12E$. Therefore 45 vibrational modes are expected, from which 32 are infrared active (B_2 and E) and 43 Raman active (A_1, B_1, B_2 and E). A Raman spectrum of solid $[\text{NEt}_4]_2[\text{Zn}(\text{dmit})_2]$ has been published.²⁴ Unfortunately, the complex decomposed in the laser beam with a 10 to 800 mW output and the use of an indium-gallium-arsenic (InGaAs) detector only led to a low resolution spectrum. In our work, the use of a germanium detector and a lower laser intensity lead to a higher resolution spectrum, without decomposition.

The infrared absorption spectrum and the Raman scattering spectrum of solid $[\text{NEt}_4]_2[\text{Zn}(\text{dmit})_2]$ are shown in Figures 2 and 3. Several main spectral features should be noted. The intense band at 1417 cm^{-1} was assigned by Cassoux³ to the C=C ligand stretching vibration. Between 1060 and 1010 cm^{-1} , two very intense bands at 1059 and 1039 cm^{-1} were assigned as C=S terminal stretching. These bands are generally used to confirm the presence of dmit ligands in complexes.^{7,8} Other assignments are: bands at 995 and 887 cm^{-1} to C-S stretches, the low intensity band at 524 cm^{-1} , a breathing mode of the heterocycle dmit

**Figure 2.** Infrared vibrational spectrum of the metal complex salt $[\text{NEt}_4]_2[\text{Zn}(\text{dmit})_2]$: (a) 4000-600 and (b) 600-150 cm^{-1} .**Figure 3.** Vibrational Raman spectrum of the metal complex salt $[\text{NEt}_4]_2[\text{Zn}(\text{dmit})_2]$: (a) 4000-600 and (b) 600-150 cm^{-1} .

complex, and the bands below 500 cm^{-1} , to angular deformation and Zn-S stretching modes.

Similar bands are found between 500 and 140 cm^{-1} in both the IR and the Raman spectra, e.g., the band at $524/525\text{ cm}^{-1}$, that is probably related to an E vibrational band. An intense band at 461 cm^{-1} with a low intensity shoulder at 450 cm^{-1} in the infrared spectrum corresponds to the single band at 465 cm^{-1} in the Raman spectrum. At 388 cm^{-1} (IR) and at 390 cm^{-1} (Raman), low intensity bands were observed. A band is observed at 349 cm^{-1} in the Raman spectrum. A central band is observed at 314 cm^{-1} in the IR spectrum, with shoulders at $322, 306$ and 300 cm^{-1} , while in the Raman spectrum, only a single band at 306 cm^{-1} is observed. The bands at 249 cm^{-1} (IR) and 247 (Raman) have not been assigned. A very strong band at 195 cm^{-1} in the IR spectrum corresponds to the low intensity Raman band at 198 cm^{-1} . Similarly, a weak IR band at 158 cm^{-1} corresponds to the strong band at 161 cm^{-1} in the Raman spectrum.

Several low intensity bands between 2100 and 1700 cm^{-1} in the IR have been assigned as overtone and combination bands, e.g., the $2065, 1986$ and 1776 cm^{-1} bands have wavenumbers twice those at $1039, 995$ and 888 cm^{-1} , respectively. Furthermore, the 2081 and 1880 cm^{-1} can be considered as the combination of $1059, 1039$ and $1417, 461\text{ cm}^{-1}$ bands.

$[Zn(dmit)_2]^{2-}$ frequency calculations

Calculations have employed two different basis sets and ECPs (SBKd and Stuttgart) and three different methodological levels (RHF, DFT and MP2) for the vibration spectra calculation. Intensities were calculated for the infrared and Raman spectra. All frequencies showed real modes as expected for an equilibrium geometry. The results are presented on Table 2, where the following conventions should apply over all article: ν = stretching mode, σ = scissoring mode, τ = twisting mode, ω = wagging mode, ρ = rocking mode, ϕ = torsion, χ = out of plane.

The most noticeable observation concerns the constancy of the results with the different calculation methods. Even in the low frequency vibrational modes related to the Zn-S stretching frequencies, the order and intensity of bands do not change either with the basis sets or with the methods. As expected, RHF calculated frequencies are greater than found experimentally, but with scaling factors smaller than those reported in the literature for most vibrations.^{33,34} An exception is the C=C stretching frequency which was calculated to be 20% greater. For the low frequency modes, the results are remarkable, with most vibrations showing differences less than 1% of the experimental results. DFT and MP2 calculations overestimated most of the stretching vibrations,

as expected, however in these cases the error was less than 7% for the DFT-B3LYP and 3% for the MP2 method. Other frequencies showed a slight tendency to be under estimated.

It is remarkable that for all calculations, the same ordering of the vibrational modes was obtained, with the exceptions for the angular deformation $\omega(SC_1S)$ out-of-plane at 515 cm^{-1} in the RHF and 445 cm^{-1} in the MP2 calculation and in the case of $\chi(SC=CS)$ out-of-plane in the C=C bond at 306 cm^{-1} in the RHF calculation. Such consistent ordering allowed the assignments of several bands not previously mentioned in the literature. Among these are the Zn-S stretching modes at 314 and 306 cm^{-1} , the Zn-S stretching coupled to the $\sigma(SC=S)$ bending at 248 cm^{-1} and the dmit bending around the zinc coordination at 195 and 158 cm^{-1} , where the first mode was clearly assigned as asymmetric bending. Modes between 350 and 320 cm^{-1} were readily assigned to the ligand in-plane bending with a small metal participation. A mode of medium intensity, not previously assigned, was the dmit ring SC_1S bend at 461 cm^{-1} , an useful mode for the dmit ring identification.

Between 1100 and 990 cm^{-1} , the experimental spectra show three intense bands assigned as C=S stretching modes. However, theoretical calculations show a pair of a_1/b_2 bands at $1057/1055\text{ cm}^{-1}$ and a pair of degenerate e bands at 1038 cm^{-1} at RHF infrared spectrum.

According to previous studies,²⁴ the assignment of these bands is not an easy task. A tentative assignment defines these three bands at $1059, 1039$ and 995 cm^{-1} as C=S stretch, C-S stretch and SC_1S angular deformation and an overtone of C-S stretch, respectively.²⁴ Other possibility would be the first as C=S stretch combined with SCS angular deformation, the second would be C=S and C-S stretch combined to the ring deformation, leaving the last unassigned.²⁴ Another possibility would consider crystalline network splitting of the e modes, not previously mentioned.

In order to explore the problem we recorded the infrared spectra in dichloromethane solution, which showed three bands with similar wavenumbers, with the central (1039 cm^{-1}) being the most intense. The presence of similar bands in solution discards any suggestion of crystalline net effect in this multiplet. Therefore, we assign this band as an asymmetrical C=S stretching while the other bands were assigned as symmetrical C=S stretching and the e C-S stretching modes. This assignment is consistent with a related unpublished work where we synthesized the corresponding dmio complex (C=S changes to C=O), observing that bands at $1039/1059\text{ cm}^{-1}$ disappear. Therefore, based on intensity arguments, our final conclusion points to an symmetrical C=S stretching at

Table 2. Comparison between the observed and calculated wavenumbers (cm⁻¹) of vibrational spectra of [NEt₄]₂[Zn(dmit)₂]

cm ⁻¹	RHF		DFT		MP2		Symm.	Exp.		Mode
	I _{IR}	I _{RA}	cm ⁻¹	I _{IR^a}	cm ⁻¹	I _{IR}		IR(s)	R(s)	
1732	0.00	213.60	1516	0	1457	0.00	A ₁			ν(C=C)
1732	2.54	11.72	1512	681	1455	13.45	B ₂	1417	1421	ν(C=C)
1057	0.00	26.14	1007	0	1021	0.00	A ₁	1059		ν(C=S)
1055	23.02	0.09	1006	738	1016	27.81	B ₂	1039	1030	ν(C=S)
1038	0.04	0.00	928	11	950	0.25	E			ν(Ce-Sm)
1038	0.04	0.00	928	11	950	0.25	E			ν(Ce-Sm)
926	0.51	10.38	830	32	836	1.10	E	910	902	ν(Ct-Sa)
926	0.51	10.39	830	32	836	1.10	E			ν(Ct-Sa)
879	0.00	7.15	803	0	819	0.00	A ₁			ν(Ce-Sm) + ν(Ce-Sa)
878	2.76	1.63	801	221	817	1.96	B ₂	887		ν(Ce-Sm) + ν(Ce-Sa)
765	0.01	0.16	700	1	714	0.10	E			ν(Ce-Sm) + ν(Ce-Sa) + ν(Ct-Sa)
765	0.01	0.15	700	1	714	0.10	E			ν(Ce-Sm) + ν(Ce-Sa) + ν(Ct-Sa)
644	0.00	4.53	528	0	512	0.00	B ₁			τ(SC=CS)
644	0.00	0.00	527	0	510	0.00	A ₂			τ(SC=CS)
537	0.00	16.38	494	0	497	0.00	A ₁			ν(C=S) + ν(Ct-S)
537	0.61	6.84	494	26	496	1.60	B ₂	524	525	ν(C=S) + ν(Ct-S)
474	0.00	6.13	439	0	438	0.00	A ₁			σ(SCTs) + ν(Ce-Sm) + ν(Ce-Sa)
474	1.75	6.00	438	89	435	3.05	B ₂	461	465	σ(SCTs) + ν(Ce-Sm) + ν(Ce-Sa)
515	0.00	0.68	436	2	445	0.10	E			ω(SCTs)
515	0.00	0.68	436	2	445	0.10	E			ω(SCTs)
411	0.06	8.63	375	4	381	0.07	E	388	390	σ(CCSm) + ρ(SCTs)
411	0.06	8.63	375	4	381	0.07	E			σ(CCSm) + ρ(SCTs)
350	0.00	0.58	329	2	333	0.25	B ₂			σ(CeSaCt) + σ(SCTs)
348	0.00	61.75	326	0	330	0.00	A ₁		349	σ(CeSaCt) + σ(SCTs)
347	0.00	1.33	321	0	327	0.00	E			σ(CeSaCt) + σ(SC=S) + ν(Zn-S)
347	0.00	1.20	321	0	327	0.00	E			σ(CeSaCt) + σ(SC=S) + ν(Zn-S)
314	3.42	4.86	305	111	316	3.01	B ₂	314		ν(Zn-S)
306	0.00	4.45	290	0	299	0.00	A ₁		306	ν(Zn-S)
307	0.00	1.74	267	0	261	0.01	E			χ(SC=CS)
307	0.00	1.74	267	0	261	0.01	E			χ(SC=CS)
249	0.49	1.05	238	14	253	0.39	E	249	247	ν(Zn-S) + σ(SC=S)
249	0.49	1.05	238	14	253	0.39	E			ν(Zn-S) + σ(SC=S)
198	1.91	0.31	192	83	198	1.56	B ₂	195	198	σ(SZnS)
158	0.21	3.93	147	4	156	0.10	E	158	161	ω(SZnS) + σ(SC=S)
158	0.21	3.91	147	4	156	0.10	E			ω(SZnS) + σ(SC=S)
131	0.00	0.06	112	0	116	0.00	B ₁			τ(SaCeSm)
130	0.00	0.00	111	0	114	0.00	A ₂			τ(SaCeSm)
105	0.17	1.28	96	5	96	0.12	E			ω(SZnS) + τ(SaCeSm)
105	0.17	1.30	96	5	96	0.12	E			ω(SZnS) + τ(SaCeSm)
95	0.00	6.79	92	0	96	0.00	A ₁			ω(SZnS)
74	0.00	3.39	70	0	74	0.00	E			σ(SZnS) + χ(SC=CS) + σ(SC=S)
74	0.00	3.39	70	0	74	0.00	E			σ(SZnS) + χ(SC=CS) + σ(SC=S)
42	0.00	9.21	32	0	39	0.00	B ₁			σ(SZnS) + χ(SmCeSa)
16	0.00	0.04	14	0	13	0.00	E			χ(SC=CS) + χ(SC=S) + σ(SZnS)
16	0.00	3.41	14	0	13	0.00	E			χ(SC=CS) + χ(SC=S) + σ(SZnS)

^a Calculated infrared intensities (I_{IR}) are in Debye²(Å⁻²·AMU). Raman scattering activities (I_{RA}) are in Å⁴AMU⁻¹. Calculated infrared intensities (I_{IR^a}) are in KMmol⁻¹. ^b Ce = carbon ethylenic, Ct = carbon of thione, Sm = linked sulfur to the metal, Sa = sulfur of the ring.

1059 cm⁻¹, an asymmetrical C=S stretching at 1039 cm⁻¹ and a C-S stretching at 995 cm⁻¹.

[Sb(dmit)₂]⁻¹ geometry optimization

The individual anion [Sb(dmit)₂]⁻¹ in [NEt₄][Sb(dmit)₂] in the solid state has C₂ symmetry. Figure 4 represents the optimised structure of the [Sb(dmit)₂]⁻¹ and its atomic numbering scheme, while Table 3 lists the structural data

obtained from the SBKd calculation using the three methods. All of which provided similar results, except for the dihedral angle SbSCC that presented a difference among the DFT method in comparison with the RHF and MP2 methods. The calculated optimised C₂ gas phase structure is just the same as that found in the solid state, ignoring the inter-anionic Sb-S interactions: the gas-phase calculated and experimental solid state intra-atomic distances differed by less than 6%.

The results between the three basis functions do not exhibit significant variations as also found for the zincate complex discussed earlier, an indication that polarization functions do not play a qualitative role in this geometry calculation. Analysing the differences between the several calculations methods, it is noticeable that the RHF method provided the closest match to the experimental values, similarly as observed for the dmit zinc complex. The description of the correlation energy through the DFT or MP2 methods showed an increase on the Sb-S bonds.

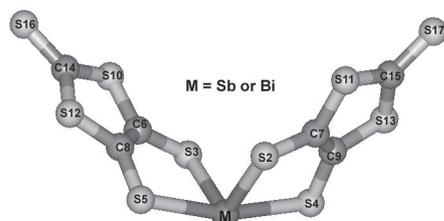


Figure 4. Three-dimensional representation of complex anions $[M(dmit)_2]^{-1}$ ($M = Sb$ or Bi) calculated with ab initio method.

Table 3. Main geometrical parameters for the complex anion $[Sb(dmit)_2]^{-1}$ with ab initio calculations and the SBKd basis set

	RHF	DFT	MP2	Exp. ⁷
<i>Bond (Å)</i>				
$S^{17}=C^{15}$	1.6809	1.6917	1.7033	1.6800
$S^{13}-C^{15}$	1.7847	1.8122	1.8237	1.7250
$S^{13}-C^9$	1.8219	1.8394	1.8420	1.7400
S^4-C^9	1.8052	1.8102	1.8243	1.7250
$C^7=C^9$	1.3226	1.3513	1.3648	1.3800
Sb-S	2.7215	2.7337	2.7345	2.6720
	2.5438	2.5808	2.5737	2.4960
<i>Bond (°)</i>				
$S^{13}C^{15}S^{17}$	123.6	124.0	123.8	123.4
$C^{15}S^{13}C^9$	96.7	97.2	96.9	97.7
$S^{13}C^9C^7$	116.9	116.7	116.9	115.6
SbS^4C^9	100.3	100.9	99.8	101.2
S^2SbS^4	82.8	83.2	83.7	83.2
S^3SbS^4	165.4	167.0	165.8	158.8
S^5SbS^2	88.0	88.9	87.4	83.5
S^3SbS^2	101.3	104.9	102.5	101.6
<i>Diedral (°)</i>				
$SbS^4C^7C^9$	13.3	8.9	14.6	8.1
$S^2C^7S^{11}C^{15}$	178.7	178.8	178.3	174.5
$C^7S^{11}C^{15}S^{17}$	178.8	178.7	178.8	174.9

$[Sb(dmit)_2]^{-1}$ experimental frequencies

The molecular cartesian displacement reducible representation for C_2 symmetry leads to the following genuine vibration space composition: $\Gamma_{\text{vibration}} = 23A + 22B$. All 45 expected vibrational modes are both infrared and Raman active.

We are not aware of any previous Raman spectra recorded for solid $[Q][Sb(dmit)_2]$ complexes. The infrared

and the Raman spectra, from 4000 to 150 cm^{-1} , for the solid $[NEt_4]_2[Sb(dmit)_2]$, are shown in Figures 5 and 6. The strong band at 1435 cm^{-1} is assigned as the C=C dmit stretch. Between 1060 and 1010 cm^{-1} two intense bands are assigned as C=S stretches. These three bands, also present in the infrared/Raman spectra of $[Q]_2[Zn(dmit)_2]$, have been used to confirm the presence of the dmit group in such complexes.¹⁰⁻¹² The bands at 993, 895 and 881 cm^{-1} are assigned as C-S stretches. At 524 cm^{-1} , the breathing mode of the dmit ligand was observed.

It should be pointed out that no assignment of vibrational frequencies below 500 cm^{-1} has been previously made, even though a list of IR spectral bands was presented in a recent article.¹⁰ The infrared spectrum between 500 and 140 cm^{-1} , showed eleven main bands, while the Raman presented just nine bands. Between 470 and 440 cm^{-1} in the infrared and Raman spectra, two bands are observed, an intense band at 461 cm^{-1} and a low intensity band at 445 cm^{-1} . A low intensity band at 386 cm^{-1} is present in the IR spectrum, but not in the Raman spectrum. There are three bands at 351, 336 and 328 cm^{-1} in the infrared spectrum: of these, the first two are Raman active as well. Between 300 and 260 cm^{-1} , three infrared active bands were observed at 295, 283 and 274 cm^{-1} , of these the first and the last were also Raman active. The peak in the IR spectrum at 220 cm^{-1} and the shoulder at 226 cm^{-1} are matched in the Raman spectra by two distinct bands at 231 and 222 cm^{-1} . At 150 cm^{-1} , in the infrared and at 158 cm^{-1} on the Raman spectra a very intense band was observed.

Several low intensity bands have been assigned in this work as overtones and combination bands. Between 2100 and 1700 cm^{-1} , not previously reported in the literature, we assign the bands at 2044, 1969 and 1757 cm^{-1} as overtones of the fundamental modes present at 1028, 993 and 881 cm^{-1} . On the other hand, the bands at 2063, 1896 and 1782 cm^{-1} correspond to the combination of the respective fundamental modes at 1049, 1028 and 1434, 463 and 895, 881 cm^{-1} . Clearly the low frequency region of the infrared and Raman spectra, from 360 to 150 cm^{-1} , can be considered as a characteristic fingerprint region for metal-dmit complexes. The larger number of bands observed in this region for $[NEt_4][Sb(dmit)_2]$ than for $[NEt_4]_2[Zn(dmit)_2]$, can be related to the reduced symmetry of the anion in $[NEt_4][Sb(dmit)_2]$ (C_2) compared to that in $[NEt_4]_2[Zn(dmit)_2]$ (D_{2d}).

$[Sb(dmit)_2]^{-1}$ frequency calculations

From the optimized geometries we carried out a hessian calculation and frequency evaluation of all vibrational modes of $[Sb(dmit)_2]^{-1}$. All frequencies showed real modes,

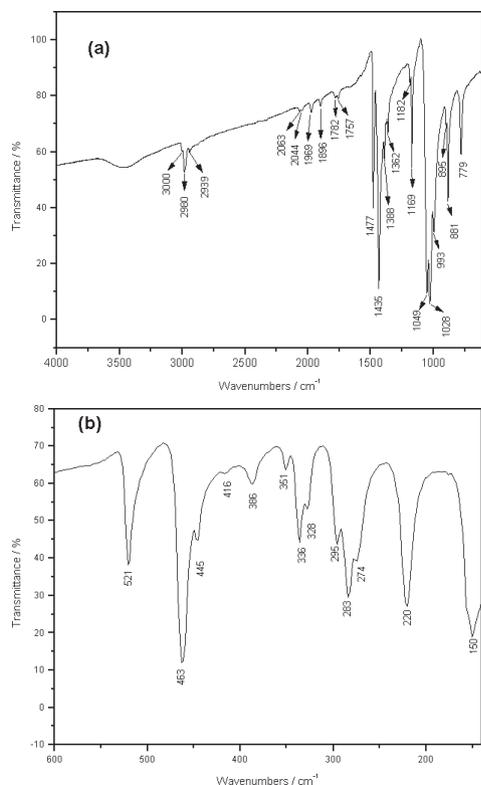


Figure 5. Infrared vibrational spectrum of the metal complex salt $[\text{NEt}_4][\text{Sb}(\text{dmit})_2]$: (a) 4000-600 and (b) 600-150 cm^{-1} .

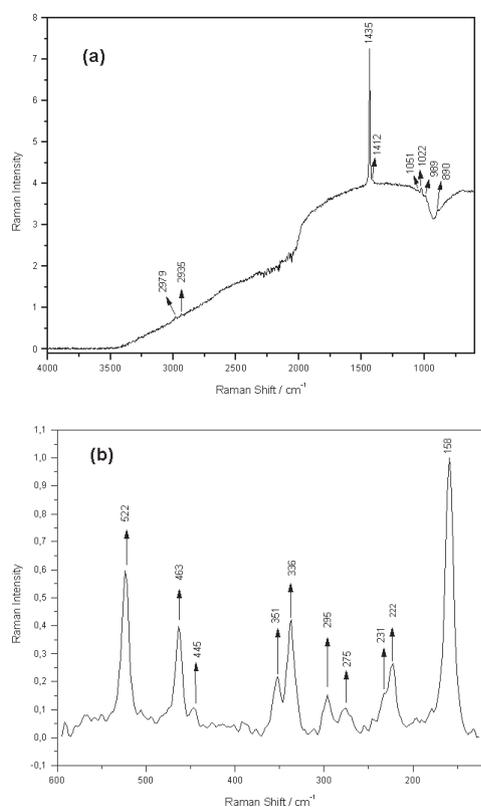


Figure 6. Vibrational Raman spectrum of the metal complex salt $[\text{NEt}_4][\text{Sb}(\text{dmit})_2]$: (a) 4000-600 and (b) 600-150 cm^{-1} .

as expected for an equilibrium geometry. The frequencies and the corresponding intensities are shown in Table 4. As observed for the zinc complex, only minor differences between the various basis sets and the metallic ECP used were detected. Larger variations were observed between the different computational methods/electron correlation descriptions. For instance, the RHF results shown in Table 4, clearly overestimated all frequencies, with the largest difference found for the C=C stretch (20%). Below 500 cm^{-1} , most of the calculated frequencies were overestimated between 5 and 10%. In the DFT and MP2 calculations, all deviations were found to be smaller, e.g., deviations for the C=C stretch were 8% and 5% from DFT-B3LYP and MP2 methods respectively. For the other modes above 300 cm^{-1} , most were overestimated between 1 to 10%. However below 300 cm^{-1} , much better agreement with experimental results were obtained, with discrepancies between 1 and 3%. Remarkable is the perfect match between the frequency ordering found among all calculations carried out and the experimental results. Two exceptions, also observed for the zinc-dmit complex, should be pointed: $w(\text{SC}_i\text{S})$ out-of-plane bending between 450 and 430 cm^{-1} , for the DFT-B3LYP and MP2 methods. This reordering was not observed for the corresponding RHF calculation between 510 and 480 cm^{-1} .

The calculations confirmed the assignments previously reported in the literature.¹⁰ The C=C, C=S and C-S stretches were recognized at 1435, 1028 and 881 cm^{-1} , respectively. The bands at 522 and 463 cm^{-1} were assigned as breathing mode from the dmit ring and SC_iS bending mode, respectively. Between 360 and 300 cm^{-1} , a series of bands was assigned as $\text{S}_a\text{C}_e\text{S}_m$ bending mode slightly coupled to the Sb-S stretching mode. The bands at 295 and 283 cm^{-1} were assigned as symmetrical and asymmetrical Sb-S stretches coupled to C=C out-of-plane twisting and $\text{SC}=\text{S}$ in-the-plane bending. These modes were assigned from their symmetry and spectroscopic activity. While the symmetric stretch is Raman-active, the asymmetrical stretch occurs in the infrared. The band at 274 cm^{-1} arises from the C=C out-of-plane. The band at 222 cm^{-1} results from the Sb-S stretching coupled to the $\text{S}=\text{C}_i\text{S}$ bending. Finally the 150 cm^{-1} band was assigned as the five-members-ring deformation formed by the metal coordination to the thiolate sulfur.

$[\text{Bi}(\text{dmit})_2]^{-1}$ geometry optimizations

The symmetry of the $[\text{Bi}(\text{dmit})_2]^{-1}$ anion in $[\text{NEt}_4][\text{Bi}(\text{dmit})_2]$ is C_2 . Table 5 lists the results for the geometry optimization carried out with the SBKd basis set and ECP pseudopotential. As previously mentioned

Table 4. Comparison between the observed and calculated wavenumbers (cm^{-1}) of vibrational spectra of $[\text{NEt}_4][\text{Sb}(\text{dmit})_2]$

cm^{-1}	RHF		DFT		MP2		Symm.	Exp.		Mode Assignment
	I_{IR}	I_{RA}	cm^{-1}	I_{IR^a}	cm^{-1}	I_{IR}		IR(s)	R(s)	
1774	0.96	17.38	1560	111	1511	2.38	B	1435	1435	$\nu(\text{C}=\text{C})$
1774	0.35	50.51	1560	34	1510	1.08	A			$\nu(\text{C}=\text{C})$
1090	10.34	0.68	1036	238	1052	8.71	A	1049		$\nu(\text{C}=\text{S})$
1088	13.88	0.48	1034	564	1050	12.39	B	1028	1022	$\nu(\text{C}=\text{S})$
1068	0.01	3.28	938	4	956	0.12	A	993	989	$\nu(\text{Ce}-\text{Sm})$
1068	0.00	4.24	937	4	956	0.08	B			$\nu(\text{Ce}-\text{Sm})$
917	0.01	23.00	819	29	832	0.64	A			$\nu(\text{Ct}-\text{Sa})$
917	0.83	3.85	816	50	830	1.22	B	895	890	$\nu(\text{Ct}-\text{Sa})$
915	0.84	21.79	816	103	814	0.09	A			$\nu(\text{Ce}-\text{Sm}) + \nu(\text{Ce}-\text{Sa})$
913	1.11	15.68	815	14	813	0.78	B	881		$\nu(\text{Ce}-\text{Sm}) + \nu(\text{Ce}-\text{Sa})$
794	0.05	5.64	712	3	718	0.09	A			$\nu(\text{Ce}-\text{Sm}) + \nu(\text{Ce}-\text{Sa}) + \nu(\text{Ct}-\text{Sa})$
793	0.20	1.02	712	12	718	0.53	B			$\nu(\text{Ce}-\text{Sm}) + \nu(\text{Ce}-\text{Sa}) + \nu(\text{Ct}-\text{Sa})$
654	0.00	0.90	540	0	529	0.00	A			$\tau(\text{SC}=\text{CS})$
652	0.01	0.73	539	0	527	0.00	B			$\tau(\text{SC}=\text{CS})$
526	0.32	22.06	482	13	476	0.87	A			$\nu(\text{C}=\text{S}) + \nu(\text{Ct}-\text{S})$
526	0.42	4.92	481	32	475	1.19	B	522	522	$\nu(\text{C}=\text{S}) + \nu(\text{Ct}-\text{S})$
487	0.59	9.26	446	11	445	0.44	A			$\sigma(\text{SCtS}) + \nu(\text{Ce}-\text{Sm}) + \nu(\text{Ce}-\text{Sa})$
486	1.24	3.92	446	23	444	0.23	B	463	463	$\sigma(\text{SCtS}) + \nu(\text{Ce}-\text{Sm}) + \nu(\text{Ce}-\text{Sa})$
509	0.00	0.57	444	0	443	0.03	A			$\omega(\text{SCtS})$
509	0.01	4.43	444	16	443	0.81	B	445	445	$\omega(\text{SCtS})$
419	0.04	7.09	380	1	384	0.03	A			$\sigma(\text{CCSm}) + \rho(\text{SCtS})$
417	0.06	5.32	379	3	382	0.05	B	386		$\sigma(\text{CCSm}) + \rho(\text{SCtS})$
379	0.12	12.83	336	8	339	0.06	A	351	351	$\sigma(\text{CeSaCt}) + \sigma(\text{SCtS}) + \nu(\text{Sb}-\text{S})$
375	0.35	2.20	335	1	338	0.24	B			$\sigma(\text{CeSaCt}) + \sigma(\text{SCtS}) + \nu(\text{Sb}-\text{S})$
360	0.02	22.18	331	1	332	0.00	A	336	336	$\sigma(\text{CeSaCt}) + \nu(\text{Sb}-\text{S})$
358	0.05	13.42	325	3	328	0.01	B	328		$\sigma(\text{CeSaCt}) + \nu(\text{Sb}-\text{S})$
330	0.28	28.42	296	8	298	0.25	A	295	295	$\nu(\text{Sb}-\text{S}) + \chi(\text{SC}=\text{CS}) + \sigma(\text{SC}=\text{S})$
326	1.31	2.92	288	109	291	2.21	B	283		$\nu(\text{Sb}-\text{S}) + \chi(\text{SC}=\text{CS}) + \sigma(\text{SC}=\text{S})$
304	0.65	9.70	270	16	265	0.08	B	274	275	$\chi(\text{SC}=\text{CS}) + \nu(\text{Sb}-\text{S})$
303	0.08	12.85	266	0	258	0.01	A			$\chi(\text{SC}=\text{CS}) + \nu(\text{Sb}-\text{S})$
255	0.00	32.88	230	0	235	0.00	A		231	$\nu(\text{Sb}-\text{S}) + \sigma(\text{SC}=\text{S})$
245	2.46	0.45	232	61	236	2.08	B	222		$\nu(\text{Sb}-\text{S}) + \sigma(\text{SC}=\text{S})$
168	2.14	10.90	152	6	158	0.17	A		158	$\sigma(\text{SSbS})$
176	0.25	10.06	158	78	158	1.84	B	150		$\sigma(\text{SSbS})$
152	1.34	1.87	142	5	146	0.10	B			$\sigma(\text{SSbS}) + \omega(\text{SC}=\text{S})$
158	0.01	8.00	131	0	143	0.01	A			$\tau(\text{SmCeSa})$
134	0.04	0.24	113	0	115	0.02	B			$\tau(\text{SmCeSa})$
125	0.02	11.69	112	0	110	0.01	A			$\tau(\text{SSbS}) + \tau(\text{SmCeSa})$
118	0.08	5.62	98	3	103	0.06	A			$\omega(\text{SSbS}) + \tau(\text{SmCeSa})$
98	0.11	2.51	92	3	89	0.07	B			$\omega(\text{SSbS}) + \tau(\text{SC}=\text{S})$
77	0.01	4.06	66	0	66	0.00	A			$\omega(\text{SSbS}) + \chi(\text{SC}=\text{CS}) + \tau(\text{SC}=\text{S})$
61	0.01	2.63	50	0	50	0.00	B			$\tau(\text{SSbS}) + \chi(\text{SmCeSa})$
53	0.00	7.72	42	0	45	0.00	A			$\chi(\text{SmCeSa})$
30	0.00	2.35	26	0	25	0.00	B			$\sigma(\text{SSbS}) + \chi(\text{SmCeSaCt})$
13	0.01	1.30	12	0	11	0.00	A			$\sigma(\text{SSbS}) + \chi(\text{SmCCSa})$

^a Calculated infrared intensities (I_{IR}) are in $\text{Debye}^2(\text{\AA}^{-2}\text{-AMU})$. Raman scattering activities (I_{RA}) are in $\text{\AA}^4\text{AMU}^{-1}$. Calculated infrared intensities (I_{IR^a}) are in KMmol^{-1} . ^b Ce = carbon ethylenic, Ct = carbon of thione, Sm = linked sulfur to the metal, Sa = sulfur of the ring.

for the zinc and antimony systems, only minor differences have been found between the calculations using the main geometrical optimized parameters with other basis sets and ECP. Therefore these results will be not presented. Similarly as found for zinc and antimony, most of the calculated geometrical parameters for the isolated anion in the gas phase showed small differences compared to those obtained in the X-ray diffraction study: differences were less than 6% for bond lengths and 10% for the angle between the

ligand planes and the dihedral angle, BiSCC. This last angle difference is considered to result from the strong inter-anionic interactions occurring in the crystalline environment, and which were neglected in the gas phase calculations. As previously noted, the RHF method provided the best match to the experimental geometrical parameters. The correlation energy description through DFT and MP2 methods improved the Bi-S bond lengths in the $[\text{NEt}_4][\text{Bi}(\text{dmit})_2]$ complex.

Table 5. Main geometrical parameters for the complex anion $[\text{Bi}(\text{dmit})_2]^{-1}$ with ab initio calculations and the SBKd basis set

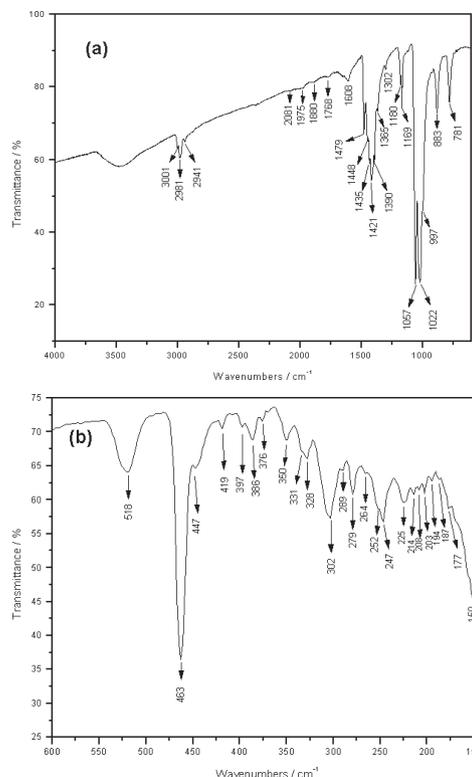
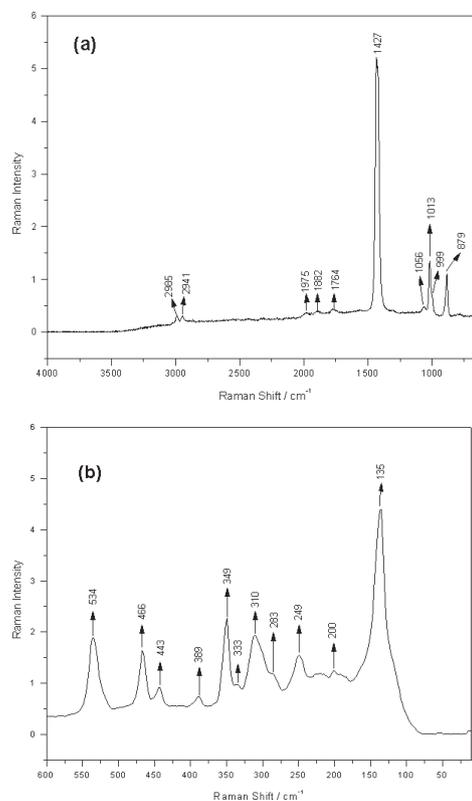
	RHF	DFT	MP2	Exp. ⁷
<i>Bond (Å)</i>				
S ¹⁷ =C ¹⁵	1.6816	1.6919	1.7035	1.6455
S ¹³ -C ¹⁵	1.7834	1.8110	1.8225	1.7231
S ¹³ -C ⁹	1.8244	1.8425	1.8452	1.7493
S ⁴ -C ⁹	1.8069	1.8099	1.8248	1.7327
C ⁷ =C ⁹	1.3233	1.3527	1.3658	1.3402
Bi-S	2.7861	2.7991	2.8053	2.8339
	2.6012	2.6565	2.6520	2.6668
<i>Bend (°)</i>				
S ¹³ C ¹⁵ S ¹⁷	123.7	124.0	123.9	124.4
C ¹⁵ S ¹³ C ⁹	96.9	97.4	97.1	99.2
S ¹³ C ⁹ C ⁷	116.8	116.6	116.7	115.1
BiS ⁴ C ⁹	100.5	100.9	100.3	96.9
S ² BiS ⁴	81.9	81.8	82.5	76.5
S ⁵ BiS ⁴	166.8	169.2	167.1	152.5
S ⁵ BiS ²	89.8	91.8	89.7	84.2
S ³ BiS ²	101.9	108.0	105.7	89.6
<i>Dihedral (°)</i>				
BiS ² C ⁷ C ⁹	9.7	5.9	9.2	30.3
S ² C ⁷ S ¹¹ C ¹⁵	179.4	179.5	179.3	177.4
C ⁷ S ¹¹ C ¹⁵ S ¹⁷	179.1	178.9	178.9	177.6

[Bi(dmit)₂]⁻¹ experimental frequencies

The $[\text{Bi}(\text{dmit})_2]^{-1}$ and $[\text{Sb}(\text{dmit})_2]^{-1}$ structures have the same symmetry and both exhibit 45 vibrational infrared and Raman active modes. The infrared and Raman spectra of solid $[\text{NEt}_4][\text{Bi}(\text{dmit})_2]$, from 4000 to 120 cm^{-1} , are shown in Figures 7 and 8.

The intense band at 1421 cm^{-1} was assigned to the C=C stretch.¹⁴ Two other very strong bands, 1057 and 1022 cm^{-1} , have been assigned as symmetrical and asymmetric C=S stretches. A band at 883 cm^{-1} and a shoulder at 997 cm^{-1} were assigned as C-S stretches. At 534 and 518 cm^{-1} , two other bands have been assigned as dmit ring breathing modes.

Six high intensity bands occur in the infrared spectrum between 500 and 140 cm^{-1} , while seven high intensity bands were detected in the Raman spectrum. In the region from 470 to 440 cm^{-1} , in both the infrared and Raman spectra, an intense band at 463 cm^{-1} and a shoulder at 447 cm^{-1} are found. Following these bands, a central and more intense band at 386 cm^{-1} and two weak bands at 397 and 376 cm^{-1} are observed. In the Raman spectrum a single band was identified at 389 cm^{-1} . Three bands were observed at 350, 331 and 328 cm^{-1} in the infrared spectrum, but only the first was actually observed in Raman spectrum. Between 320 and 300 cm^{-1} , a shoulder at 310 cm^{-1} and a band at 302 cm^{-1} are observed in infrared spectrum. In the Raman spectrum, the main band in this region was observed at 310 cm^{-1} with a shoulder at 302 cm^{-1} . At 279 cm^{-1} in the

**Figure 7.** Infrared vibrational spectrum of the metal complex salt $[\text{NEt}_4][\text{Bi}(\text{dmit})_2]$: (a) 4000-600 and (b) 600-150 cm^{-1} .**Figure 8.** Vibrational Raman spectrum of the metal complex salt $[\text{NEt}_4][\text{Bi}(\text{dmit})_2]$: (a) 4000-600 and (b) 600-10 cm^{-1} .

infrared spectrum a band was observed, while the Raman band in this region was found at 283 cm⁻¹ with a low intensity. In the region between 260 and 240 cm⁻¹, two bands at 252 and 247 cm⁻¹ were observed in the infrared, while in the Raman spectrum just one band was observed at 249 cm⁻¹. Between 230 and 180 cm⁻¹, several low intensity bands were observed in the infrared and Raman. At 134 cm⁻¹ in the Raman spectrum, a single intense band

was observed, assigned as the five-members-ring deformation formed by the metal coordination to the thiolate sulfur. This band has been found to have a linear dependency with the mass of the M-S bond along the bis(dmit) complex of zinc, antimony and bismuth.

The region from 2100 to 1700 cm⁻¹ was investigated for overtones and combination bands. The bands at 2081 and 1880 cm⁻¹, might be interpreted as combination of the

Table 6. Comparison between the observed and calculated wavenumbers (cm⁻¹) of vibrational spectra of [NEt₄][Bi(dmit)₂]

cm ⁻¹	RHF		DFT		MP2		Symm.	Exp.		Mode Assignment
	I _{IR*}	I _{RA}	cm ⁻¹	I _{IR*}	cm ⁻¹	I _{IR}		IR(s)	R(s)	
1770	43	17	1553	93	1506	103	B	1427	1427	$\nu(\text{C}=\text{C})$
1770	14	63	1553	25	1505	40	A			$\nu(\text{C}=\text{C})$
1089	374	0	1037	201	1051	295	A	1057	1056	$\nu(\text{C}=\text{S})$
1086	684	0	1035	641	1049	675	B	1020	1013	$\nu(\text{C}=\text{S})$
1059	1	3	934	4	950	5	A	997	999	$\nu(\text{Ce}-\text{Sm})$
1059	0	3	933	3	950	5	B			$\nu(\text{Ce}-\text{Sm})$
918	0	24	818	6	822	24	A			$\nu(\text{Ct}-\text{Sa})$
918	21	6	817	37	821	70	B			$\nu(\text{Ct}-\text{Sa})$
904	30	39	809	29	816	3	A		879	$\nu(\text{Ce}-\text{Sm}) + \nu(\text{Ce}-\text{Sa})$
903	69	16	808	124	814	35	B	883		$\nu(\text{Ce}-\text{Sm}) + \nu(\text{Ce}-\text{Sa})$
791	2	6	711	2	717	4	A			$\nu(\text{Ce}-\text{Sm}) + \nu(\text{Ce}-\text{Sa}) + \nu(\text{Ct}-\text{Sa})$
790	8	1	711	11	716	20	B			$\nu(\text{Ce}-\text{Sm}) + \nu(\text{Ce}-\text{Sa}) + \nu(\text{Ct}-\text{Sa})$
652	0	1	540	0	528	0	A			$\tau(\text{SC}=\text{CS})$
650	0	1	539	0	526	0	B			$\tau(\text{SC}=\text{CS})$
528	11	23	483	10	477	30	A	534	534	$\nu(\text{C}=\text{S}) + \nu(\text{Ct}-\text{S})$
528	20	5	482	34	476	64	B	518		$\nu(\text{C}=\text{S}) + \nu(\text{Ct}-\text{S})$
485	22	10	445	9	443	6	A			$\sigma(\text{SCtS}) + \nu(\text{Ce}-\text{Sm}) + \nu(\text{Ce}-\text{Sa})$
484	59	4	445	27	441	54	B	463	466	$\sigma(\text{SCtS}) + \nu(\text{Ce}-\text{Sm}) + \nu(\text{Ce}-\text{Sa})$
510	1	5	444	0	444	13	A	447	443	$\omega(\text{SCtS})$
510	0	0	443	13	444	2	B			$\omega(\text{SCtS})$
415	1	8	377	1	382	2	A	386	389	$\sigma(\text{CCSm}) + \rho(\text{SCtS})$
414	2	5	376	3	382	4	B			$\sigma(\text{CCSm}) + \rho(\text{SCtS})$
371	2	8	332	0	337	11	B	350	349	$\sigma(\text{SmCeSa}) + \sigma(\text{SC}=\text{S}) + \nu(\text{Bi}-\text{S})$
367	8	2	333	4	344	14	A			$\sigma(\text{SmCeSa}) + \sigma(\text{SC}=\text{S}) + \nu(\text{Bi}-\text{S})$
358	0	17	329	0	330	0	A	331	333	$\sigma(\text{CeSaCt}) + \sigma(\text{CtSaCe}) + \nu(\text{Bi}-\text{S})$
354	1	11	322	2	329	0	B	328		$\sigma(\text{CeSaCt}) + \sigma(\text{CtSaCe}) + \nu(\text{Bi}-\text{S})$
324	12	35	293	6	298	4	A		310	$\nu(\text{Bi}-\text{S}) + \chi(\text{SC}=\text{CS}) + \sigma(\text{SC}=\text{S})$
320	46	2	282	103	287	108	B	302		$\nu(\text{Bi}-\text{S}) + \chi(\text{SC}=\text{CS}) + \sigma(\text{SC}=\text{S})$
300	65	11	270	40	264	5	B	279	283	$\chi(\text{SC}=\text{CS}) + \nu(\text{Bi}-\text{S})$
304	4	17	270	0	262	0	A			$\chi(\text{SC}=\text{CS}) + \nu(\text{Bi}-\text{S})$
256	0	27	232	0	240	0	A			$\nu(\text{Bi}-\text{S}) + \sigma(\text{SC}=\text{S})$
243	93	1	229	40	234	75	B	247	249	$\nu(\text{Bi}-\text{S}) + \sigma(\text{SC}=\text{S})$
152	3	1	134	16	143	3	B			$\sigma(\text{SBiS})$
148	9	4	133	3	139	12	A		135	$\sigma(\text{SBiS}) + \sigma(\text{SC}=\text{S})$
138	96	6	133	50	133	52	B			$\sigma(\text{SBiS}) + \tau(\text{SmCeSa})$
142	6	12	119	2	126	2	A			$\sigma(\text{SBiS}) + \sigma(\text{SC}=\text{S})$
122	0	8	111	0	109	0	A			$\sigma(\text{SBiS}) + \tau(\text{SmCeSa})$
128	2	0	110	0	108	0	B			$\tau(\text{SBiS}) + \tau(\text{SmCeSa})$
108	5	15	93	3	95	4	A			$\omega(\text{SBiS}) + \tau(\text{SmCeSa}) + \chi(\text{SC}=\text{CS}) + \tau(\text{SC}=\text{S})$
95	5	3	90	3	86	3	B			$\omega(\text{SBiS}) + \tau(\text{SC}=\text{S})$
71	0	3	59	0	59	0	A			$\omega(\text{SBiS}) + \chi(\text{SC}=\text{CS}) + \tau(\text{SC}=\text{S})$
53	1	2	44	0	45	0	B			$\chi(\text{SmCCSa}) + \sigma(\text{SBiS}) + \tau(\text{SC}=\text{S})$
47	0	8	37	0	43	0	A			$\chi(\text{SmCeSa}) + \sigma(\text{SBiS})$
25	0	3	24	0	23	0	B			$\sigma(\text{SBiS}) + \chi(\text{SmCeSaCt})$
10	0	12	11	0	10	0	A			$\sigma(\text{SBiS}) + \chi(\text{SmCCSa})$

^a Calculated infrared intensities (I_{IR*}) are in K²mol⁻¹. Raman scattering activities (I_{RA}) are in Å⁴AMU⁻¹. ^b Ce = carbon ethylenic, Ct = carbon of thione, Sm = linked sulfur to the metal, Sa = sulfur of the ring.

1057, 1022 cm^{-1} and 1421, 463 cm^{-1} bands. On the other hand, the bands at 1975 and 1764 have been assigned as overtones of the 997 and 883 cm^{-1} fundamental bands.

[Bi(dmit)₂]⁻¹ frequency calculations

From the optimized geometries a hessian matrix calculation and frequency determination of all vibrational modes of $[\text{Bi}(\text{dmit})_2]^{-1}$ was carried out. All frequencies showed real modes. The frequencies and the corresponding intensities are shown in Table 6. As observed previously, only minor differences between the various basis sets/metal ECPs were detected. The results followed trends similar to the ones previously discussed in the zinc and antimony cases.

Conclusions

This work reports a combined theoretical-experimental investigation of the infrared and Raman vibrational spectra of three metal complexes with the 1,3-dithiole-2-thione-4,5-dithiolate (dmit) ligand, namely, $[\text{NEt}_4]_2[\text{Zn}(\text{dmit})_2]$, $[\text{NEt}_4][\text{Sb}(\text{dmit})_2]$ and $[\text{NEt}_4][\text{Bi}(\text{dmit})_2]$. The IR and Raman spectra of all complexes have been recorded, and assigned, with an extensive set of *ab initio* calculations carried out for validation and cross-checking.

Calculations carried out with several basis set and ECP potentials generated very similar results among them, as well as to the X-ray diffraction experimental results. Actually only minor differences were pointed in the text. The authors, however, claim this should not be considered a general conclusion for other metals and complexes. The metal-ligand character of dmit-complexes might be an important part of this observed behaviour. For other metal complexes and metal clusters molecular systems an additional study must be carried out. Within this possibility we still expect that the ECP and basis set choice should have an important role in the quality of the results. Major differences have been observed however with the use of different methodologies. RHF showed the worst results while DFT and MP2 calculations presented equivalent quality.

Starting from the observation of the spectra in solution of the $[\text{NEt}_4]_2[\text{Zn}(\text{dmit})_2]$, C=S symmetrical and asymmetrical stretching modes were defined more clearly than in the solid state, becoming easily distinguishable and coherent with the present theoretical observations in 1059 and 1039 cm^{-1} .

Most of the spectral data exhibit a spectrum composed of isolated fundamental bands for the systems without a strong M-S inter-molecular interaction. Within this pattern very few overtones, combinations bands and crystalline

lattice vibrations have been observed in the spectra of the dmit-metal complexes. For the zinc complex, independently of the cation size, a similar spectrum was observed. On the other hand, for the dmit compounds where strong M-S inter-molecular interactions are present, significant deviations are observed, depending on the cation size and crystalline environment. The most clarifying example was found for the tetraethylammonium bismuth dmit complex for which a large difference among the BiSCC dihedral angle (20/25 degrees) was found between calculated and experimental results. Similarly significant overtones, combinations bands and crystalline effects appear in the spectrum. Although present on bismuth, for the antimony complex, whose is the same but with weaker M-S inter-molecular interactions, the spectra showed a simple fundamental behaviour.

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Electronic Supplementary Information

Theoretical and vibrational analysis of tetraethylammonium bromide are available as a PDF file at <http://jbcs.sbq.org.br>

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Vibrational Spectra of Bis(dmit) Complexes of Main Group Metals: IR, Raman and *Ab Initio* Calculations

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Results

Tetraethylammonium bromide NEt_4Br

The infrared and Raman spectra of NEt_4Br were obtained and a complete assignment of all the vibrations associated with the cation was made, aided by *ab initio* calculations.

$[NEt_4]^+$ geometry optimization

Several symmetries have been found for the $[NEt_4]^+$ cation³⁵⁻³⁷ in its various salts: the one found by X-ray crystallography for the salts studied here is D_{2d} ,^{7,12,15} Figure 1. As no metal atom is present, all atoms in the calculations were represented by a single basis set. Table 1 lists the main geometrical parameters calculated by the different *ab initio* methods for the $[NEt_4]^+$ cation. It is clear that the calculated main geometric parameters are insensitive to the method used, all indeed giving geometric parameters close to those found in the the X-ray diffraction study.

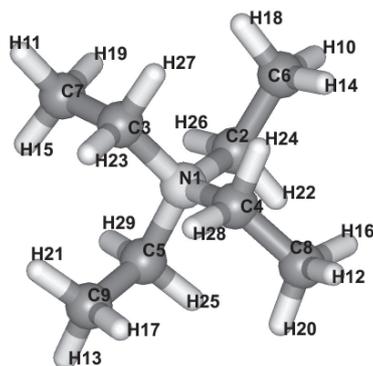


Figure 1. Three-dimensional representation of tetraethylammonium calculated with *ab initio* method.

Table 1. Main geometrical parameters for the tetraethylammonium cation calculated with *ab initio* methods and 6-31G basis set

	RHF	DFT	MP2	Exp. ⁷
<i>Bond</i> (Å)				
N ¹ -C ²	1.5287	1.5469	1.5518	1.5176
C ² -C ⁶	1.5223	1.5255	1.5345	1.5031
C ² -H ¹⁶	1.0780	1.0916	1.0956	0.9699
C ⁶ -H ¹⁰	1.0830	1.0934	1.0986	0.9601
<i>Bend</i> (°)				
C ⁵ NC ²	106.2	106.5	106.4	106.9
C ⁶ C ² N	116.0	115.8	115.5	115.6
C ² C ⁶ H ¹⁴	113.1	113.1	112.7	109.5
C ² C ⁶ H ¹⁰	107.5	107.6	107.4	109.5
C ² C ⁶ H ¹⁸	113.1	113.1	112.7	109.5
NC ² H ²⁶	106.1	105.7	105.6	108.4
NC ² H ²²	106.1	105.7	105.6	108.3
<i>Dihedral</i> (°)				
NC ² C ⁶ H ¹⁰	180.0	180.0	180.0	176.4
NC ² C ⁶ H ¹⁴	62.2	62.1	62.1	56.4
NC ² C ⁶ H ¹⁸	62.2	62.2	62.1	63.6
H ¹⁰ C ⁶ C ² H ²²	59.7	59.7	59.9	65.1
H ¹⁰ C ⁶ C ² H ²⁶	59.3	59.7	59.9	58.6

$[NEt_4]^+$ experimental frequencies

For the D_{2d} point group, the molecular cartesian displacement representation and its irreducible composition lead to 81 vibration modes; $\Gamma_{\text{vibration}} = 12A_1 + 8A_2 + 9B_1 + 12B_2 + 20E$. From these, 52 are infrared active (B_2 and E) and 73 are Raman active (A_1 , B_1 , B_2 and E).

The infrared and Raman spectra of several tetraethylammonium salts have been reported.^{31,32} Figures 2 and 3 show the infrared and Raman spectra of tetraethylammonium bromide, respectively. According to Hooper and James,³¹ the modes at 2999 and 2983 cm^{-1} in the infrared spectrum are assigned as asymmetric stretches, within CH_3 and CH_2 respectively, while the bands at 2970 and 2944 cm^{-1} in the Raman spectrum are symmetrical CH_3 and CH_2 stretches, respectively. Bands between 2960

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and 1500 cm^{-1} are, most probably, overtone and combinations bands. The region between 1500 and 1100 cm^{-1} is characteristic of HCH bending modes, including scissoring, wagging, twisting and rocking modes. A set of bands between 1493 and 1440 cm^{-1} has been assigned as asymmetric and symmetric CH_3 scissoring modes. The bands between 1415 and 1390 cm^{-1} and the band at 1378 cm^{-1} are assigned as CH_2 scissoring modes. The Raman active band at 1304 cm^{-1} has been assigned as CH_2 twisting mode. Two very intense bands at 1183 and 1175 cm^{-1} in the infrared spectrum have been assigned to CH_3 asymmetrical rocking modes and that at 1116 cm^{-1} to a symmetrical rocking mode in the Raman spectrum.

In the region below 1100 cm^{-1} , many C-N and C-C stretching modes are present. Low intensity bands at 1071 and 1050 cm^{-1} have been assigned as C-C asymmetrical stretching. A shoulder at 1030 cm^{-1} has been assigned as an asymmetric C-N stretch, while the intense band present at 1009 cm^{-1} in the infrared spectrum is assigned as a CH_3 rocking mode. Two low intensity bands at 910 and 898 cm^{-1} have been considered C-C symmetric stretching, while the bands at 802 and 789 cm^{-1} have been assigned as rocking of the CH_2 group. The Raman active band at 679 cm^{-1} was considered by Hooper and James as the C-N

characteristic symmetric stretching mode.

In the low wavenumber region of the spectra, the few low intensity bands have been assigned to the deformation of the CCN frame. Only in the Raman spectrum was observed an intense mode at 422 cm^{-1} . The 476 , 421 and 390 cm^{-1} bands in the infrared spectrum have been assigned to the asymmetric CNC, CCN and symmetric CNC vibrational modes. Hooper and James do not inform about the presence of bands below 390 cm^{-1} .

$[\text{NEt}_4]^+$ frequency calculations

Frequencies and intensities were calculated for the infrared and Raman spectra. All frequencies showed real modes, as expected for an equilibrium geometry. The proposed assignments, based on *ab initio* calculations for all 81 vibrational modes present in NEt_4Br , is found on Table 2, with the previous conventions adopted.

For this molecule we employed RHF, DFT and MP2 methods with a single basis set to check the identification of the IR/Raman cationic bands in the spectra of the complex salts.^{33,34} The most significant bands of $[\text{NEt}_4]^+$ were observed from 3000 to 2900 cm^{-1} , from 1500 to 1100 cm^{-1} , at 785 cm^{-1} and at 419 cm^{-1} .

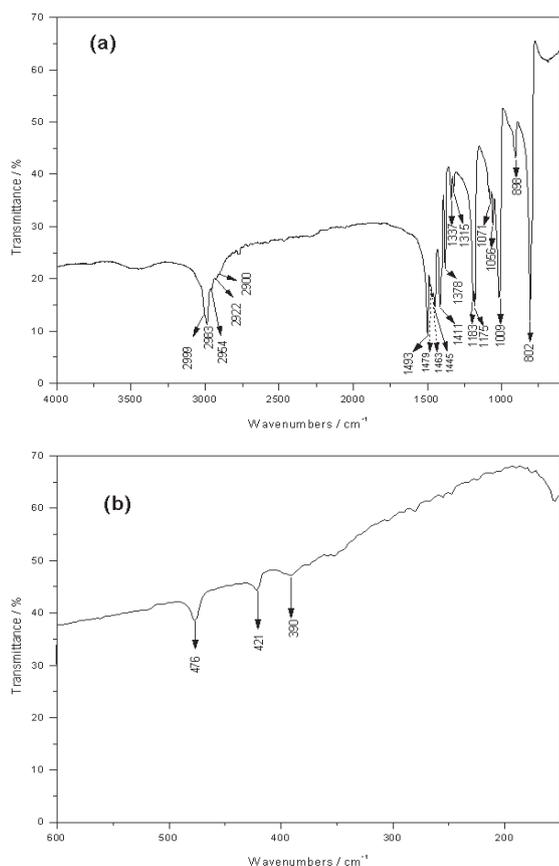


Figure 2. Infrared vibrational spectrum of the tetraethylammonium bromide: (a) 4000 - 600 and (b) 600 - 100 cm^{-1} .

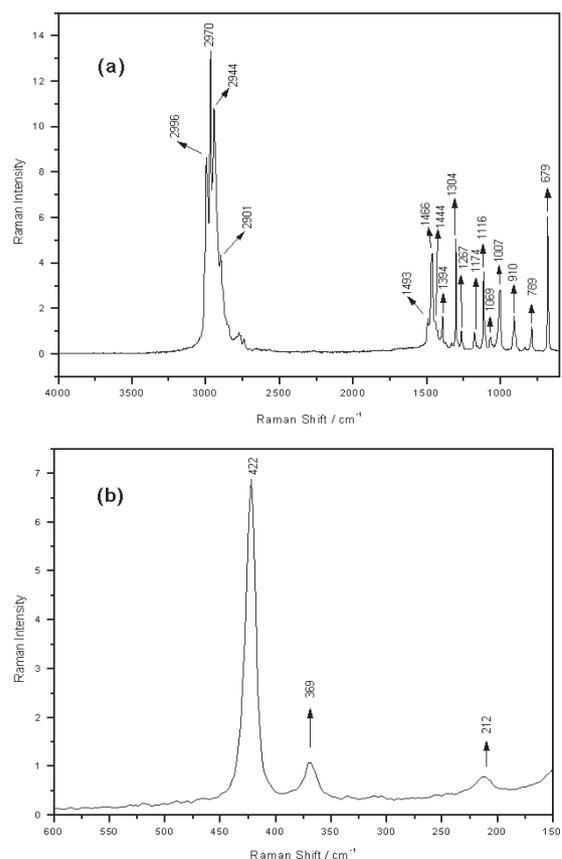


Figure 3. Vibrational Raman spectrum of the tetraethylammonium bromide: (a) 4000 - 600 and (b) 600 - 150 cm^{-1} .

Table 2. Comparison between the observed and calculated wavenumbers (cm^{-1}) of vibrational spectra of NEt_4Br

cm^{-1}	RHF		DFT		MP2		Symm.	Exp.		Mode Assignment
	I_{IR}	I_{RA}	cm^{-1}	I_{IR^*}	cm^{-1}	I_{IR}		IR(s)	R(s)	
3346	0.50	42.2	3180	26	3193	0.65	E			$\nu(\text{C-H}) / (\text{CH}_3+\text{CH}_2)$
3346	0.50	42.2	3180	26	3193	0.65	E	2999	2996	$\nu(\text{C-H}) / (\text{CH}_3+\text{CH}_2)$
3342	0.00	3.0	3179	0	3195	0.00	B_1			$\nu(\text{C-H}) / (\text{CH}_3+\text{CH}_2)$
3331	0.00	0.0	3169	0	3182	0.00	A_2			$\nu(\text{C-H}) / (\text{CH}_3+\text{CH}_2)$
3307	0.38	120.3	3155	6	3175	0.06	E			$\nu(\text{C-H}) / (\text{CH}_3+\text{CH}_2)$
3307	0.38	120.3	3155	6	3175	0.06	E			$\nu(\text{C-H}) / (\text{CH}_3+\text{CH}_2)$
3306	0.00	0.0	3156	0	3174	0.00	A_2			$\nu(\text{C-H}) / (\text{CH}_3+\text{CH}_2)$
3301	1.12	15.3	3149	28	3172	0.44	B_2	2983		$\nu(\text{C-H}) / (\text{CH}_3+\text{CH}_2)$
3298	0.00	33.0	3147	0	3170	0.00	B_1			$\nu(\text{C-H}) / (\text{CH}_3+\text{CH}_2)$
3292	0.00	31.4	3145	0	3168	0.00	E			$\nu(\text{C-H}) / (\text{CH}_2)$
3292	0.00	31.4	3145	0	3168	0.00	E			$\nu(\text{C-H}) / (\text{CH}_2)$
3292	0.00	166.3	3146	0	3163	0.00	A_1		2970	$\nu(\text{C-H}) / (\text{CH}_3+\text{CH}_2)$
3279	0.00	160.3	3119	10	3123	0.36	B_2	2954		$\nu(\text{C-H}) / (\text{CH}_2)$
3289	0.00	96.0	3121	0	3122	0.00	A_1		2944	$\nu(\text{C-H}) / (\text{CH}_2)$
3275	0.11	24.4	3112	3	3115	0.05	E	2922		$\nu(\text{C-H}) / (\text{CH}_2)$
3275	0.11	24.4	3112	3	3115	0.05	E			$\nu(\text{C-H}) / (\text{CH}_2)$
3221	0.12	2.7	3075	6	3087	0.10	B_2			$\nu(\text{C-H}) / (\text{CH}_3)$
3220	0.18	15.6	3075	3	3087	0.06	E	2900		$\nu(\text{C-H}) / (\text{CH}_3)$
3220	0.18	15.6	3075	3	3087	0.06	E			$\nu(\text{C-H}) / (\text{CH}_3)$
3220	0.00	392.2	3074	0	3086	0.00	A_1			$\nu(\text{C-H}) / (\text{CH}_3)$
1693	1.49	0.2	1570	52	1590	1.23	B_2	1493	1493	$\sigma(\text{HCH}) / (\text{CH}_2)$
1691	0.00	9.3	1571	0	1588	0.00	A_1			$\sigma(\text{HCH}) / (\text{CH}_2)$
1679	0.19	1.0	1566	17	1585	0.36	E	1479		$\sigma(\text{HCH}) / (\text{CH}_2+\text{CH}_3)$
1679	0.19	1.0	1566	17	1585	0.36	E			$\sigma(\text{HCH}) / (\text{CH}_2+\text{CH}_3)$
1663	0.00	59.4	1547	0	1574	0.00	A_1			$\sigma(\text{HCH}) / (\text{CH}_3)$
1657	0.03	5.0	1546	4	1570	0.06	E			$\sigma(\text{HCH}) / (\text{CH}_3+\text{CH}_2)$
1657	0.03	5.0	1546	3	1570	0.06	E			$\sigma(\text{HCH}) / (\text{CH}_3+\text{CH}_2)$
1657	0.00	80.8	1549	0	1569	0.00	B_1		1466	$\sigma(\text{HCH}) / (\text{CH}_3)$
1656	0.00	0.0	1548	0	1568	0.00	A_2			$\sigma(\text{HCH}) / (\text{CH}_3)$
1650	0.51	0.0	1537	33	1563	0.54	B_2	1463		$\sigma(\text{HCH}) / (\text{CH}_3)$
1642	0.77	0.3	1524	22	1547	0.48	E	1445	1444	$\sigma(\text{HCH}) / (\text{CH}_3+\text{CH}_2)$
1642	0.77	0.3	1524	22	1547	0.48	E			$\sigma(\text{HCH}) / (\text{CH}_3+\text{CH}_2)$
1601	0.00	13.9	1482	0	1500	0.00	A_1		1394	$\omega(\text{HCH}) / (\text{CH}_3)$
1604	0.04	2.6	1481	3	1500	0.07	B_2			$\omega(\text{HCH}) / (\text{CH}_2)$
1600	0.81	1.2	1480	35	1500	0.78	E	1411		$\omega(\text{HCH}) / (\text{CH}_3)$
1600	0.81	1.2	1480	35	1500	0.78	E			$\omega(\text{HCH}) / (\text{CH}_3)$
1587	0.00	5.6	1461	0	1474	0.03	B_2	1378		$\omega(\text{HCH}) / (\text{CH}_2+\text{CH}_3)$
1582	0.00	4.3	1456	0	1469	0.00	A_1			$\omega(\text{HCH}) / (\text{CH}_2+\text{CH}_2)$
1533	0.61	0.8	1409	18	1421	0.32	E	1335		$\tau(\text{HCH}) + \omega(\text{HCH}) / (\text{CH}_2)$
1533	0.61	0.8	1409	18	1421	0.32	E			$\tau(\text{HCH}) + \omega(\text{HCH}) / (\text{CH}_2)$
1506	0.02	0.7	1388	0	1396	0.04	E	1317		$\tau(\text{HCH}) + \omega(\text{HCH}) / (\text{CH}_2)$
1506	0.02	0.7	1388	0	1396	0.04	E			$\tau(\text{HCH}) + \omega(\text{HCH}) / (\text{CH}_2)$
1455	0.00	36.0	1343	0	1351	0.00	B_1		1304	$\omega(\text{HCH}) / (\text{CH}_2)$
1433	0.00	0.0	1324	0	1337	0.00	A_2			$\omega(\text{HCH}) / (\text{CH}_2)$
1352	0.11	0.1	1242	2	1266	0.04	E			$\rho(\text{HCH}) / (\text{CH}_2)$
1352	0.11	0.1	1242	2	1266	0.04	E			$\rho(\text{HCH}) / (\text{CH}_3)$
1306	0.00	0.0	1203	0	1222	0.00	B_1			$\rho(\text{HCH}) / (\text{CH}_2+\text{CH}_3)$
1306	1.31	4.1	1194	29	1212	0.66	B_2	1185/75	1174	$\omega(\text{HCH}) / (\text{CH}_3)$
1248	0.00	13.3	1157	0	1168	0.00	A_1		1116	$\omega(\text{HCH}) / (\text{CH}_2)$
1244	0.00	0.0	1140	0	1157	0.00	A_2			$\rho(\text{HCH}) + \sigma\text{CCN}(\text{CH}_3 + \text{CH}_2)$
1200	0.19	2.5	1109	7	1116	0.13	E	1071	1069	$\omega(\text{HCH}) / (\text{CH}_3)$
1200	0.19	2.5	1109	7	1116	0.13	E			$\omega(\text{HCH}) / (\text{CH}_2)$
1141	0.50	1.5	1065	16	1082	0.33	B_2	1056		$\nu(\text{C-C}) + \nu(\text{N-C})$
1121	0.00	4.4	1052	0	1058	0.00	A_1			$\nu(\text{C-C}) + \nu(\text{N-C})$
1094	0.64	9.8	1020	21	1030.6	0.42	E	1010	1007	$\nu(\text{C-C}) + \nu(\text{N-C})$
1094	0.64	9.8	1020	21	1031	0.42	E			$\nu(\text{C-C}) + \nu(\text{N-C})$
990	0.08	8.4	904	0	910	0.00	B_2	898	910	$\omega(\text{CCH}) + \sigma(\text{CCN})$
974	0.03	1.2	902	0	909	0.00	E			$\nu(\text{N-C}) + \rho(\text{CCH})$
974	0.03	1.2	902	0	909	0.00	E			$\nu(\text{N-C}) + \rho(\text{CCH})$
916	0.00	0.7	850	0	852	0.00	B_1			$\rho(\text{HCH}) / (\text{CH}_2+\text{CH}_3)$
862	0.00	0.0	801	0	796	0.00	A_2			$\rho(\text{HCH}) / (\text{CH}_2+\text{CH}_3)$

Table 2. Comparison between the observed and calculated wavenumbers (cm⁻¹) of vibrational spectra of NEt₄Br

cm ⁻¹	RHF		DFT		MP2		Symm.	Exp.		Mode Assignment
	I _{IR}	I _{RA}	cm ⁻¹	I _{IR*}	cm ⁻¹	I _{IR}		IR(s)	R(s)	
846	1.22	3.4	777	43	782	0.76	E	802	789	$\nu(\text{N-C}) + \rho(\text{HCH})(\text{CH}_2+\text{CH}_3)$
846	1.22	3.4	777	43	782	0.76	E			$\nu(\text{N-C}) + \rho(\text{HCH})(\text{CH}_2+\text{CH}_3)$
711	0.00	12.3	662	0	654	0.00	A ₁		679	$\nu(\text{N-C}) + \sigma(\text{CCN}) + \omega(\text{HCH})$
564	0.00	0.2	532	0	525	0.01	E			$\sigma(\text{CCN}) + \rho(\text{HCH})$
564	0.00	0.2	532	0	525	0.01	E			$\sigma(\text{CCN}) + \rho(\text{HCH})$
441	0.03	5.2	414	1	420	0.03	B ₂	476		$\sigma(\text{CCN})$
433	0.00	9.9	406	0	407	0.00	A ₁	421	422	$\sigma(\text{CCN})$
394	0.00	0.0	369	0	370	0.00	E	390		$\rho(\text{CNC})$
394	0.00	0.0	369	0	370	0.00	E			$\rho(\text{CNC})$
377	0.00	0.6	356	0	360	0.00	B ₁		370	$\tau(\text{CNC})$
364	0.00	0.0	352	0	337	0.00	A ₂			$\phi(\text{C-C})$
311	0.00	0.1	298	0	288	0.00	B ₂			$\rho(\text{CCN}) + \rho(\text{HCH})$
254	0.00	0.0	270	0	237	0.00	B ₁			$\phi(\text{CH}_3)$
254	0.00	0.1	267	0	237	0.00	E			$\phi(\text{CH}_3)$
254	0.00	0.1	267	0	237	0.00	E			$\phi(\text{CH}_3)$
204	0.00	0.2	196	0	191	0.00	A ₁		212	$\sigma(\text{CCN})$
121	0.00	0.0	129	0	114	0.00	A ₂			$\phi(\text{C-C})$
110	0.02	0.0	106	1	103	0.01	E			$\sigma(\text{CCN})$
110	0.02	0.0	106	1	103	0.01	E			$\sigma(\text{CCN})$
75	0.00	0.2	75	0	59	0.00	B ₁			$\phi(\text{C-N})$

^a Calculated infrared intensities (I_{IR}) are in Debye²(Å⁻²-AMU). Raman scattering activities (I_{RA}) are in Å⁴AMU⁻¹. Calculated infrared intensities (I_{IR*}) are in KMmol⁻¹.