Charged Higgs Production at Photon Colliders in 2HDM-III

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We study charged Higgs production in the process $\gamma\gamma \to A^0 \to W^- H^+$. The processes $\gamma\gamma \to A^0$ are loop mediated in a 2HDM. This is due to the fact that photons only couple directly to charged particles and the Higgs only couples to particles with mass acquired via Higgs mechanism. Although in MSSM the contribution from the process $\gamma\gamma \to A^0$ is too small, it has been found that in a more general 2HDM it could be enhanced. On the other hand, the boson A^0 can decay in $W^- H^+$ at tree level and the charged Higgs can decay in fermions. So, the whole process under study is $\gamma\gamma \to A^0 \to (W^- \to l\nu)(H^+ \to f_i f_j)$ in 2HDM-III. Evidence about charged Higgs existence could demonstrate that structure of the Higgs sector has several multiplets.

Keywords: Two Higgs Doublet Model; Photon Collider; 2HDM; Higgs

I. INTRODUCTION

The Standard Model of particle physics (SM) has been succesful in explaining experimental data so far. However, the Higgs sector of the model is experimentally unknown. This fact has lead many theorists to suggest that the sector could be non-minimal, and today it is common to study extensions of the SM scalar sector. The simplest extension is the Two Higgs Doublet Model (2HDM)[1, 2], which involves two scalar doublets in the process of electroweak symmetry breaking. After the symmetry breaking, there are five physical Higgs particles: two charged Higgs $H^\pm$, two CP-even H^0, h^0, and one CP-odd A^0. The charged particles $H^\pm$ are characteristic of the models with two Higgs doublets and its discovery would be a clear signal of physics beyond the SM.

A promising alternative for the search of new physics lies on the collisions of photons that will be present at ILC (International Linear Collider). The ILC is the next greatest project to be developed after the LHC (Large Hadron Collider). The photons will be produced by Compton retrodispersion, and earlier studies show that production of neutral scalars by photon collisions present a considerable probability of detection. Besides, photons couple directly to charged particles, so $\gamma\gamma$ high energy collisions could provide a better understanding of several aspects of the SM and its extensions[4].

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In this work we study the process $\gamma\gamma \to A^0 \to (W^- \to l\nu)(H^+ \to f_i f_j)$ in the frame of a general 2HDM. The first section presents a brief overview of the 2HDM, focusing in the third type (2HDM-III). The next section contains the expressions used in the calculations. The third section contains the obtained results and finally the conclusions are drawn.

II. THE TWO HIGGS DOUBLET MODEL TYPE III

The minimal extension of the Higgs sector of the SM consists in adding a second scalar doublet with the same quantum numbers than the first one[2, 3]. We denote them by:

$$\Phi_1 = \begin{pmatrix} \phi_1^+ \\ \phi_1^0 \end{pmatrix} \qquad \Phi_2 = \begin{pmatrix} \phi_2^+ \\ \phi_2^0 \end{pmatrix}, \qquad (1)$$

$$\langle \Phi_1 \rangle_0 = \frac{v_1}{\sqrt{2}}, \qquad \langle \Phi_2 \rangle_0 = \frac{v_2}{\sqrt{2}}. \qquad (2)$$

The most general gauge invariant lagrangian which couples both Higgs fields with fermions is:

$$\begin{aligned} -\mathcal{L}_Y &= \eta_{ij}^{U,0} \bar{Q}_{iL}^0 \tilde{\Phi}_1 U_{jR}^0 + \eta_{ij}^{D,0} \bar{Q}_{iL}^0 \Phi_1 D_{jR}^0 \\ &+ \xi_{ij}^{U,0} \bar{Q}_{iL}^0 \tilde{\Phi}_2 U_{jR}^0 + \xi_{ij}^{D,0} \bar{Q}_{iL}^0 \Phi_2 D_{jR}^0 \\ &+ \eta_{ij}^{E,0} \bar{l}_{iL}^0 \tilde{\Phi}_1 E_{jR}^0 + \xi_{ij}^{E,0} \bar{l}_{iL}^0 \Phi_2 E_{jR}^0 + h.c. \end{aligned} \qquad (3)$$

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where $\eta^{U,D}$ and $\xi^{U,D}$ are non-diagonal mixing matrices 3×3, $\tilde{\Phi}_i = i\sigma_2\Phi_i$, $(U,D)_R$ are right-handed fermion singlets, Q_L are left-handed fermion doublets, and the index 0 indicates that the fields are not mass eigenstates.

In the most general case, both Higgs doublets couple with the up and down sectors, and therefore they contribute simultaneously in the proccess of mass generation for quarks. This case leads to FCNC (Flavor Changing Neutral Currents) at tree level, because it is impossible to diagonalize simultaneously both matrices η and ξ. This general case is known as 2HDM type III. However, FCNC proccesses at tree level are highly supressed by the experiment. In order to avoid their existence, Glashow and Weinberg[3] designed the following set of discrete symmetries:

$$\Phi_1 \to \Phi_1 \text{ and } \Phi_2 \to -\Phi_2,$$
$$D_{jR} \to \mp D_{jR} \text{ and } U_{jR} \to -U_{jR}. \qquad (4)$$

The condition of invariance under this discrete symmetry leads to two cases:

- By using $D_{jR} \to -D_{jR}$, matrices $\eta^{U,0}$ and $\eta^{D,0}$ have to be eliminated from the lagrangian. In this case Φ_1 decouples in the Yukawa sector and only Φ_2 gives masses to sectors up and down. This case is known as 2HDM type I.

- By using $D_{jR} \to D_{jR}$, matrices $\eta^{U,0}$ and $\xi^{D,0}$ must be eliminated from the lagrangian. In this case Φ_1 couples to the down sector and Φ_2 gives masses to up sector. This case is known as 2HDM type II.

The 2HDM-III is the only 2HDM that allows FCNC proccesses at tree level. Precission tests of the SM model show a great aggreement with the FCNC parameters, except for the phenomenon of neutrino oscillation. Besides, the FCNC processes don't seem to violate any fundamental law. We study the 2HDM-III because it has a more rich phenomenology, and it is possible to find the first two types as limit cases of this one.

In the 2HDM-III, a rotation of the scalar fields does not change the physical content of the model[3]. This rotation can get rid of the VEV of one doublet. If we take $\langle\Phi_2\rangle = 0$, it is found that $\tan\beta = 0$. This is known as the **fundamental parameterization**. We denote the VEV of the first doublet as $\langle\Phi_1\rangle = v$.

For a better study of the FCNC processes, Cheng, Sher and Yuang (CSY) propose an anzats for the Yukawa matrices such that

$$\xi_{ij} = \frac{\sqrt{m_i m_j}}{v}\lambda_{ij} \qquad (5)$$

This anzats obeys to the fact that couplings between fermions and the Higgs particle in the SM are proportional to the mass of the fermion. Parameters λ_{ij} could change the hierarchy between fermionic couplings and because of this it is expected that they would be ~ 1. Restrictions over parameters have been obtained in references [3, 5, 8]. The most relevant are:

$$
\begin{aligned}
\xi_{\mu\tau}^2 &\in [7.62 \times 10^{-4}; 4.44 \times 10^{-2}] \\
\xi_{\tau\tau} &\in [-1.8 \times 10^{-2}; 2.2 \times 10^{-2}] \\
\xi_{\mu\mu} &\in [-0.12; 0.12] \\
\xi_{\mu e} &\in [-0.39; 0.39] \\
\lambda_{bb} &\in [-6; 6] \\
\lambda_{tt} &\in [-\sqrt{8}; \sqrt{8}].
\end{aligned} \qquad (6)
$$

III. THE PROCESS $\gamma\gamma \to A^0 \to W^+ H^- \to l\nu f_i f_j$

Loops contributing in neutral Higgs production are shown in Fig. 1 and the decay $A^0 \to H^\pm W^\mp$ exists at tree level in the framework of the 2HDM-III. The process $H^- \to q_i\bar{q}_j$ has been studied in 2HDM-III[7], and under the restrictions mentioned above, it has been found that the most relevant decay is $H^- \to t\bar{b}$.
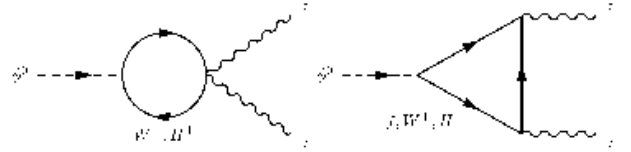

FIG. 1: Contributing diagrams in the process $\gamma\gamma \to A^0$

The decay width of the process $\gamma\gamma \to A^0$ is given by

$$\Gamma\left(\gamma\gamma \to A^0\right) = \frac{\alpha^2 g^2}{1024\pi^3} \frac{m_{A^0}^3}{m_W^2} \left| \sum_i N_C e_i^2 F(\tau) R_i^{A^0} \right|^2. \qquad (7)$$

The factor $R_i^{A^0}$ is the relative coupling between the 2HDM-III and the SM, the kinematical factor is $\tau = 4m_i^2/m_{A^0}^2$, the function $F(\tau)$ is defined as:

$$F(\tau) = -2\tau(1 + (1-\tau)f(\tau)), \qquad (8)$$

and the function f is:

$$f(\tau) = \begin{cases} -\frac{1}{4}\left|\mathrm{Ln}\left(\frac{1+\sqrt{1-\tau}}{1-\sqrt{1-\tau}}\right) - i\pi\right|^2 & \tau < 1 \\ \mathrm{ArcSin}\left(\sqrt{\frac{1}{\tau}}\right)^2 & \tau \geq 1 \end{cases}. \qquad (9)$$

Figure 2 shows the decay width of the process $\gamma\gamma \to A^0$ for several values of the parameters λ_{bb} and λ_{tt} according to the
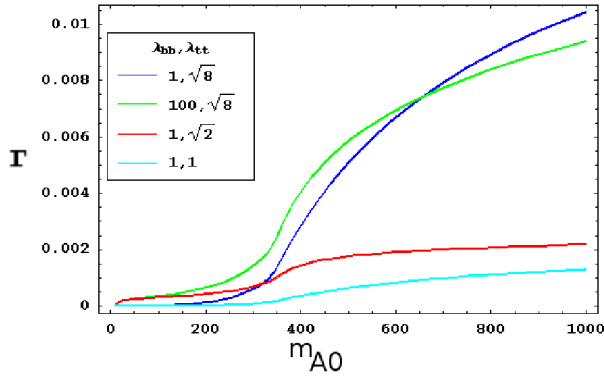
FIG. 2: Decay width of the process $\gamma\gamma \to A^0$ to one loop for $\lambda_{bb} = 1$ and $\lambda_{tt} = \sqrt{8}$

restrictions mentioned in equation 6. It is found that changes in λ_{bb} do not have a big impact on the results, while changes in λ_{tt} are considerable. It is found also that the decay width increases with the A^0 mass. So we will consider $\lambda_{bb} = 6$ and values for $m_{A^0} > 600$GeV.

The decay width for the process $A^0 \to W^- H^+$ is given by

$$\Gamma(A^0 \to W^- H^+) = \frac{\cos^2 \alpha m_{H^+}^3}{64\pi v^2} \times$$

$$\left[\sqrt{1 - \left(\frac{m_{W^-} + m_{A^0}}{m_{H^+}}\right)^2} \sqrt{1 - \left(\frac{m_{W^-} - m_{A^0}}{m_{H^+}}\right)^2} \right]^3, \quad (10)$$

where α is the mixing angle of the Higgs eigenstates, v is the vacuum expectation value as defined in the fundamental parameterization.

It has been found that the most relevant decay of the charged Higgs into fermions for this model is the decay into $t\bar{b}$ quarks[7]. This is given by

$$\Gamma(H^+ \to t\bar{b}) = \frac{3m_{H^+}^2 K_{tb}^2}{16\pi v^2} \times$$

$$\left[\left(1 - \frac{m_t^2 + m_b^2}{m_{H^+}^2}\right) - 4\lambda_{tt}\lambda_{bb}\frac{m_t^2 m_b^2}{m_{H^+}^2} \right] \times |\vec{p}_H(m_t, m_b)|, \quad (11)$$

where K_{tb} is the CKM matrix element, and $|\vec{p}_H(1,2)|$ is defined as:

$$|\vec{p}_H(m_1, m_2)| =$$

$$\sqrt{1 - \left(\frac{m_1 + m_2}{m_H}\right)^2} \sqrt{1 - \left(\frac{m_1 - m_2}{m_H}\right)^2}. \quad (12)$$

The cross section for the whole process is calculated as

$$\sigma(\gamma\gamma \to A^0 \to \tau\nu_\tau t\bar{b}) = \quad (13)$$

$$8\pi \frac{\Gamma(\gamma\gamma \to A^0)\Gamma(A^0 \to \tau\nu_\tau t\bar{b})}{(E_{\gamma\gamma}^2 - m_{A^0}^2)^2 + \Gamma_{A^0}^2 m_{A^0}^2} g(\lambda\lambda')$$

Figures 3 and 4 show the cross section for the process $\gamma\gamma \to A^0 \to W^+ H^- \to (\tau\nu_\tau)(t\bar{b})$, using $E_{\gamma\gamma} = 1000$GeV.

IV. CONCLUSSIONS
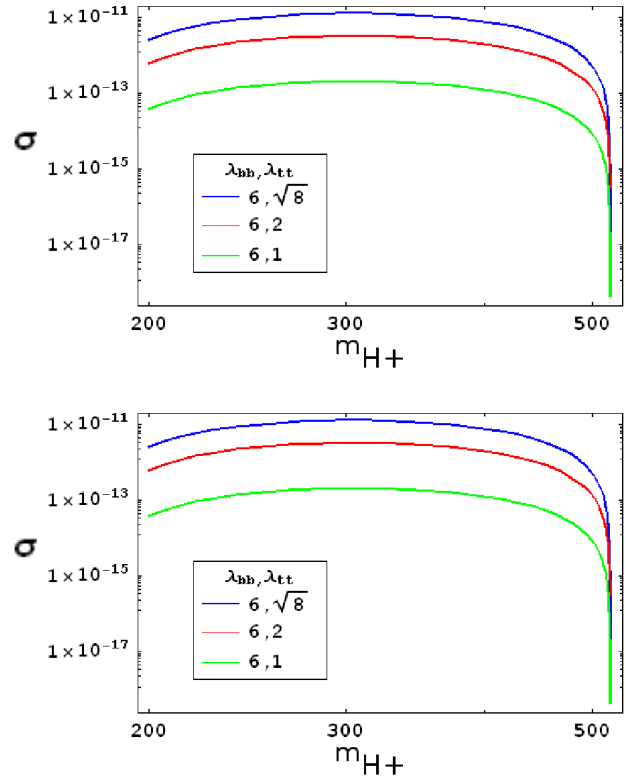

FIG. 3: Cross section (in barns) of the process $\gamma\gamma \to A^0 \to W^- H^+ \to e\nu_e t\bar{b}$ for different values of λ_{tt}, using $m_{A^0} = 600$GeV and $\lambda_{bb} = 6$

We found the cross section for the process $\gamma\gamma \to A^0 \to W^+ H^- \to (\tau\nu_\tau)(t\bar{b})$ in the frame of the 2HDM-III. Results are shown in Figs. 3 and 4. It is found that the cross section takes values between 1pb and 10pb for the parameters $\lambda_{bb} = 6$, $\lambda_{tt} = \sqrt{8}$ and $m_{A^0} = 600$GeV. For a higher value of m_{A^0} we get higher values of the cross section, between 10pb and 100pb. For lower values of the parameter λ_{tt}, we still get cross section values between 0.1pb and 10pb.

Finally, we can say that for two Higgs doublets models type III, the contribution of the process $\gamma\gamma \to A^0 \to W^- H^+ \to e\nu_e t\bar{b}$ is important, even though it is loop mediated. Earlier studies show that the contribution of this kind of processes in models such as MSSM is null[9]. The presence of this process would
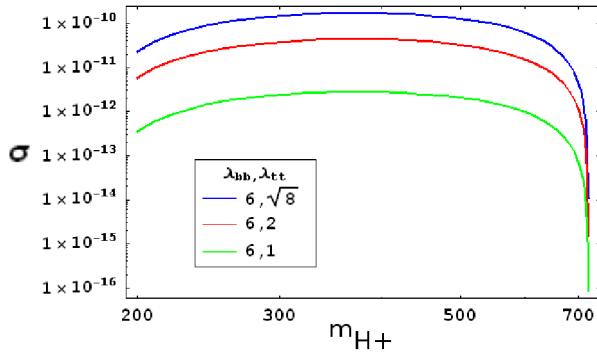
help to diferentiate between the MSSM and a more general 2HDM. Besides, evidence of the charged Higgs existence would demonstrate the multiple doublets structure of the Higgs sector.

FIG. 4: Cross section (in barns) of the process $\gamma\gamma \to A^0 \to W^- H^+ \to e\nu_e t\bar{b}$ for different values of λ_{tt}, using $m_{A^0} = 800\text{GeV}$ and $\lambda_{bb} = 6$

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