SHORT-CUTS ON WAVE REFRACTION COMPUTATION

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Synopsis

It is shown how to get rid of iterative process to compute wavelengths and how the celerities and their derivatives can be computed in terms of wavelengths for any depth, without using hyperbolic functions.

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1 - Background

It is well known that the pressure wave records have a strong attenuation of the high frequency components of the wave train, and that such records can not be corrected in the time domain. In fact the Corrections would be feasible in the frequency domain, after the Fourier analysis of the pressure records, in terms of the frequency of the Fourier lines. Such corrections are usually computed iteratively because they are given in terms of the wavelength which is the independent variable in the frequency equation. This problem was solved using polynomials and we can show that a similar solution can be found to compute celerities.

According to the wave linear theory the celerity corresponding to a wave number k and a depth d is given by

$$c = \sqrt{\frac{g}{k} \tanh (kd)}. \tag{1a}$$

In addition for deep seas

$$c_0^2 = g/k_0$$

where

$$k_0 = 2\pi/L_0 \tag{1b}$$

L_0 being the deep-sea wave length. Thus, expression (1a) can be changed into

$$\frac{c^2}{c_0^2} \frac{k}{k_0} = \tanh (kd).$$

Now

$$c = L/T$$

and

$$c_0 = L_0/T$$

where T is the period considered as being constant; consequently,

$$\frac{L^2 k}{L_0^2 k_0} = \tanh (kd).$$

Expression (1b) for $L_0 = L$ and $k_0 = k$ divided by (1b) gives

$$L/L_0 = k_0/k,$$

hence the previous equation can be transformed into

$$\frac{k_0}{k} = \tanh(kd)$$

thus

$$k_0 d = kd \tanh(kd) \tag{1c}$$

or if we make

$$k_0 d = z_0$$

and

$$kd = \underline{z} \tag{1d}$$

we obtain

$$z_0 = z \tanh z. \tag{1e}$$

The main problem is now to solve this equation for z. Since the series development of tanh z converges for

$z < \pi/2$ which corresponds to a relative depth $d/L \leq 1/4$, Elias (1971) had the idea of working out a series development to use it up to $d/L \leq 1/4$.

However if we plot the equation (1e) we see that for $z > 1.36$ ($d/L > 0.21$) the curve is nearly a straight line (Fig. 1). In the present paper, the author proposes to fitting a polynomial to the section of the curve between the abscissae 1.36 and 5.00. Several polynomials were tried and it was concluded that a 4th degree polynomial was the adequate one.

If the values of z and z_0 are known it is easy to obtain any function of these elements. Value of z obtained from iterative process and those computed with the two polynomials proved to be very close. The same technique can then be used in the computation of wave refraction were the celerities and their derivatives can be expressed in terms of z and z_0 only.
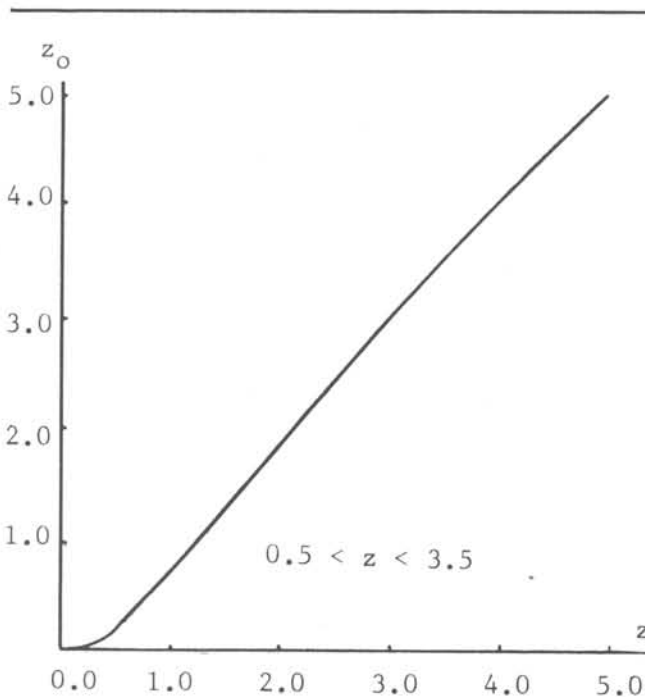

Fig. 1. Function $z_0 = z \tanh z$.

2 - Finding z

Let us develop (1e) $z_0 = z_0(z)$ according to the powers of z:

$$z_0 = z(z + a_1 z^3 + a_2 z^5 + a_3 z^7) \qquad (2a)$$

where

$$a_1 = -1/3, \quad a_2 = 2/15 \quad \text{and} \quad a_3 = -17/315 \qquad (2b)$$

We derive from (2a)

$$z^2 = z_0/(1 + a_1 z^2 + a_2 z^4 + a_1 z^6)$$

$$\simeq z_0(1 - \varepsilon + \varepsilon^2 - \varepsilon^3)$$

where

$$\varepsilon = a_1 z^2 + a_2 z^4 + a_2 z^6$$

If we use this value of ε in the preceding expression and neglect the powers of z higher than the 6th we obtain:

$$z^2 = z_0 \Big[1 - a_1 z^2 + (a_1^2 - a_2) z^4 -$$

$$(a_3 + 2a_1 a_2 + a_1^3) z^6 \Big]$$

or

$$z^2 = z_0(1 + A_1 z^2 + A_2 z^4 + A_3 z^6) \qquad (2c)$$

where

$$A_1 = -a_1 = 1/3$$
$$A_2 = a_1^2 - a_2 = -1/45 \qquad (2d)$$

and

$$A_3 = -(a_1 + 2a_1 a_2 + a_1^3) = 2/945$$

Let us now solve (2c) for z. If we assume that the solution must be compatible with (2c) we can write

$$z^2 = z_0(1 + B_1 z_0 + B_2 z_0^2 + B_3 z_0^3) \qquad (2e)$$

where B_i can be found as follows. If we use this value of z^2 in the right hand side of (2c) and neglect powers of z_0 larger than the fourth we have

$$z^2 = z_0 \Big[1 + A_1 z_0(1 + B_1 z_0 + B_2 z_0^2) + A_2 z_0^2 \times$$

$$\times (1 + 2B_1 z_0) - A_3 z_0^3 \Big] =$$

$$= z_0 \Big[1 + A_1 z_0 + (A_1 B_1 + A_2) z_0^2 +$$

$$(A_1 B_2 + 2A_2 B_1 + A_2) z_0^3 \Big] \qquad (2f)$$

Equating the coefficients of the equal powers of z_0 in (2e) and (2f) and solving for B_i we find

$$B_1 = A_1 = 1/3$$

$$B_2 = A_1 B_1 + A_2 = 4/45$$

$$B_3 = A_1 B_2 + 2A_2 B_1 + A_3 = 16/945 \qquad (2g)$$

These are the coefficients of (2e), which are very accurate if z_0 does not exceed 1.35772. For $1.35772 \leq z \leq 4.99954$ we tried polynomials from the second to the fourth degree, the latter being the best fitted:

$$z = 0.53753 + 0.54931\, z_0 + 0.14440\, z_0^2 - $$

$$0.02085\, z_0^3 + 0.00114\, z_0^4 \qquad (2h)$$

For $z_0 \geq 5$ we can assume that $z = z_0$

3 - Finding the celerity and its derivatives

Since $\omega = ck$ expression (1a) can be changed into:

$$c = \frac{g}{ck}\ \tanh z = f(z) \qquad (3a)$$

or according to (1e)

$$c = \frac{gz_0}{\omega z} = f(z) \qquad (3b)$$

Using the subscript (i,j) to denote the derivatives with respect to x or y and (ij) the derivative with respect to x and y we obtain from (3a)

$$c_{i,j} = f'(z)z_{i,j} \qquad (3c)$$

$$c_{ij} = f'(z)z_{ij} + f''(z)z_i z_j \qquad (3d)$$

where

$$f'(z) = \frac{g}{\omega}\ \mathrm{sech}^2 z \qquad (3e)$$

and

$$f''(z) = -\frac{g}{\omega}\ \mathrm{sech}^2 z\ \tanh z \qquad (3f)$$

Using (1e) to express z_i, z_j and z_{ij} in terms of $(z_0)_i$, $(z_0)_j$ and $(z_0)_{ij}$ we have after algebraic strait forward manipulations:

$$c_{i,j} = \frac{g}{\omega}\ \frac{1}{A+z}\ (z_0)_{i,j} \qquad (3g)$$

and

$$c_{ij} = \frac{g}{\omega}\ \{\frac{1}{A+z}(z_0)_{ij} -$$

$$\frac{2A(z+zz_0+Az_0-z_0)}{z_0(A+z)^3} \times (z_0)_i(z_0)_j\} \qquad (3h)$$

where

$$A = \frac{1}{2}\ \sinh 2z = \tanh z/(1-\tanh^2 z)$$

or, according to (1b),

$$A = zz_0/(z^2 - z_0^2) \qquad (3i)$$

But we have from (1a)

$$(z_0)_{i,j} = k_0 d_{i,j}$$

and

$$(z_0)_{ij} = k_0 d_{ij}$$

hence we can change (3g) and (3h), respectively, into

$$c_{i,j} = \frac{k_0 g}{\omega}\ \frac{1}{A+z}\ d_{i,j} \qquad (3j)$$

and

$$c_{ij} = \frac{k_0 g}{\omega}\ \{\frac{1}{A+z}\ d_{ij} -$$

$$\frac{2k_0 A [z+z_0(A+z)-z_0]}{z_0(A+z)^3}d_i d_j\} \qquad (3k)$$

4 - Conclusion

The Wilson (1966)'s method to compute wave refraction can be improved by introducing the above formulation for computing wavelengths, celerities and derivatives of wave celerities. If so, the computer's time is reduced by a factor of two.

References

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